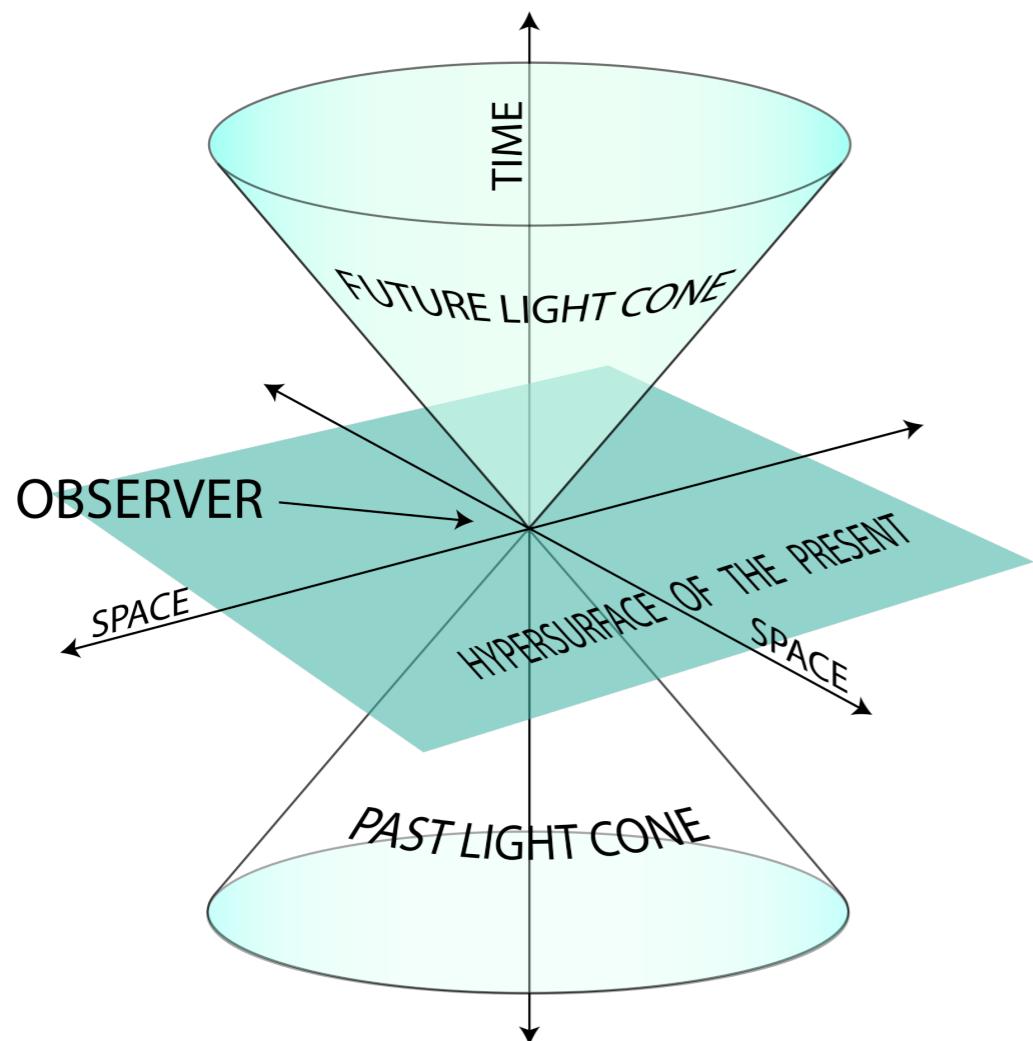


BSA of VMP in (3+1) scalar field model



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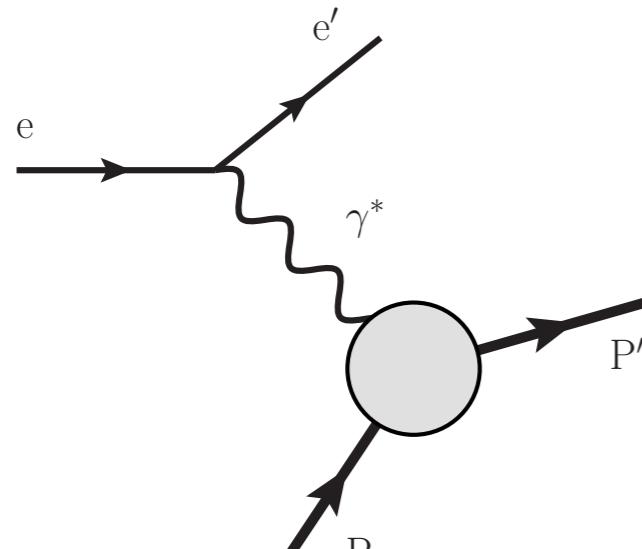
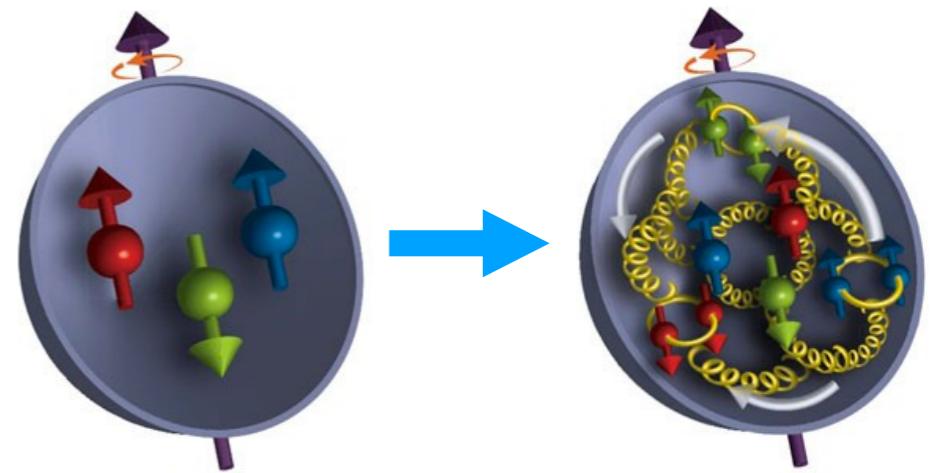


- I. Aim of this work
- II. One-Loop Scalar Model
- III. Deeply Virtual Limit
- IV. Compare Two Model

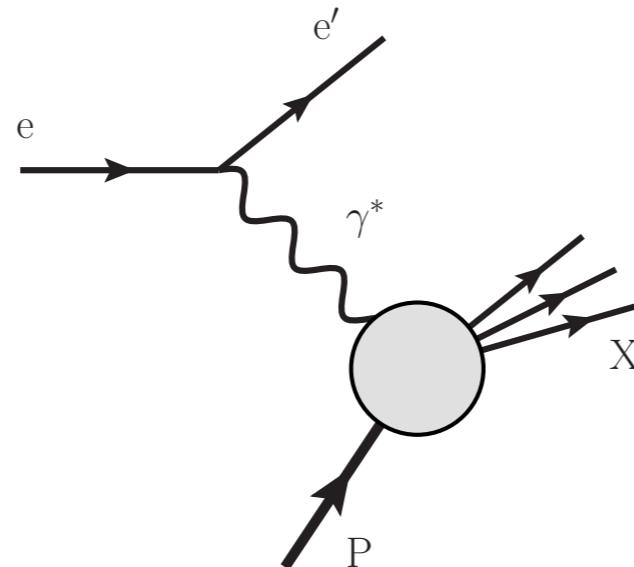
I. Aim of this work

Hadron Structure

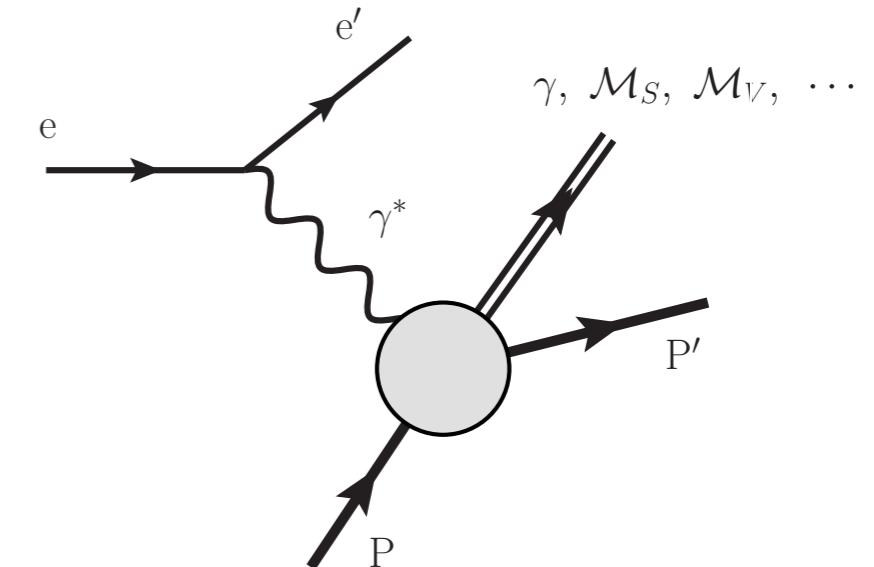
Hadron (meson and baryon) is a composite system made of **quarks**, **anti-quarks**, and **gluons**, held together by **non-perturbative QCD interaction**.



Elastic scattering



Inclusive / Inelastic



Exclusive / Inelastic

Density

$$\langle P' | \bar{\psi}(0) \hat{\mathcal{O}} \psi(0) | P \rangle$$

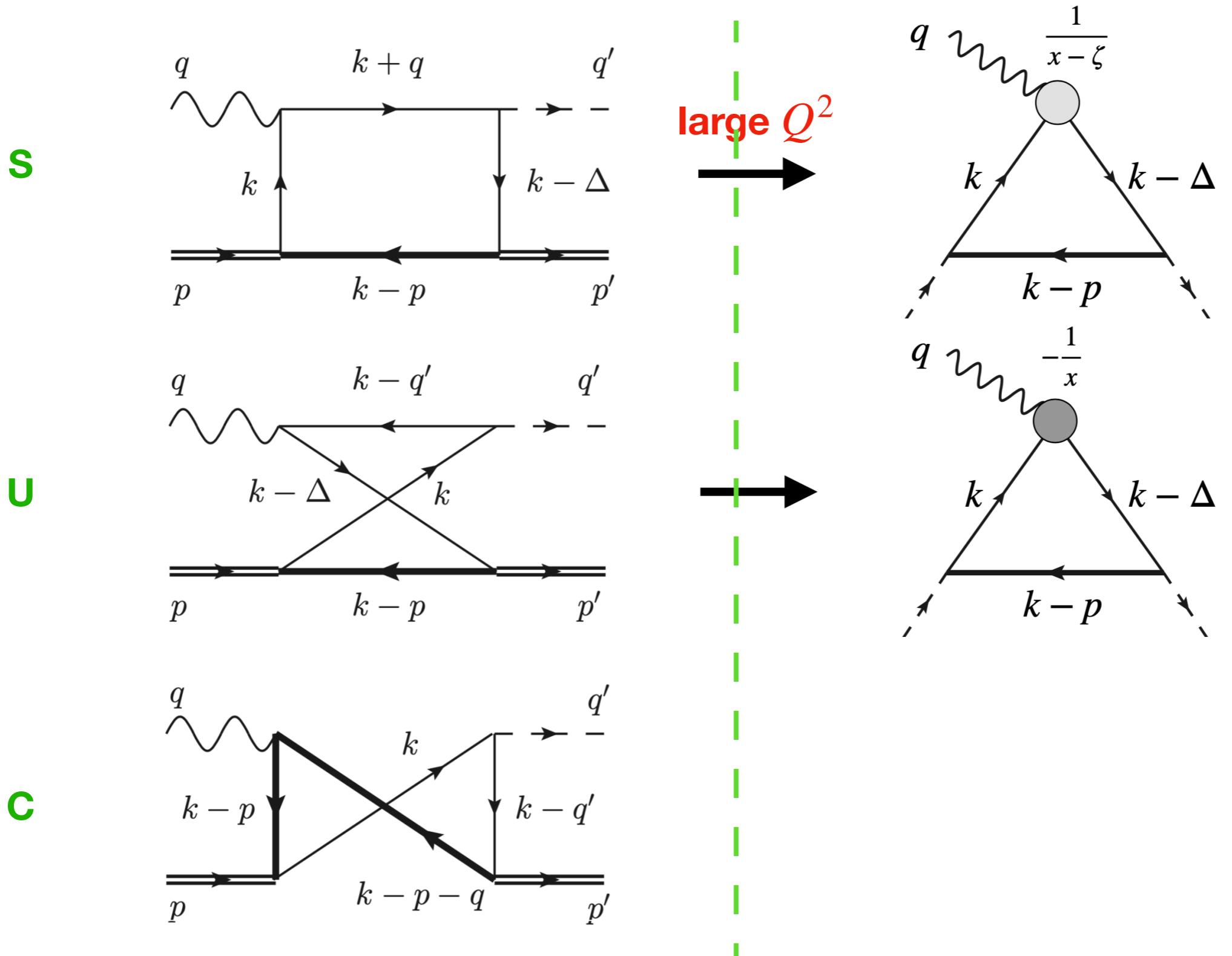
Momentum

$$\langle P | \bar{\psi}(0) \hat{\mathcal{O}} \psi(y) | P \rangle$$

Angular momentum

$$\langle P' | \bar{\psi}(0) \hat{\mathcal{O}} \psi(y) | P \rangle$$

One-loop scalar model and deeply virtual limit



One-loop scalar model

Deeply virtual limit

Aim of this work

Virtual Meson Production

$$(\gamma^* + {}^4\text{He} \rightarrow f^0 + {}^4\text{He})$$

One-loop Scalar Model
(Higher twist cont.)

Deeply Virtual Limit
(Leading twist)

Two Compton Form Factors

GPDs & PDFs & GFFs

Beam Spin Asymmetry

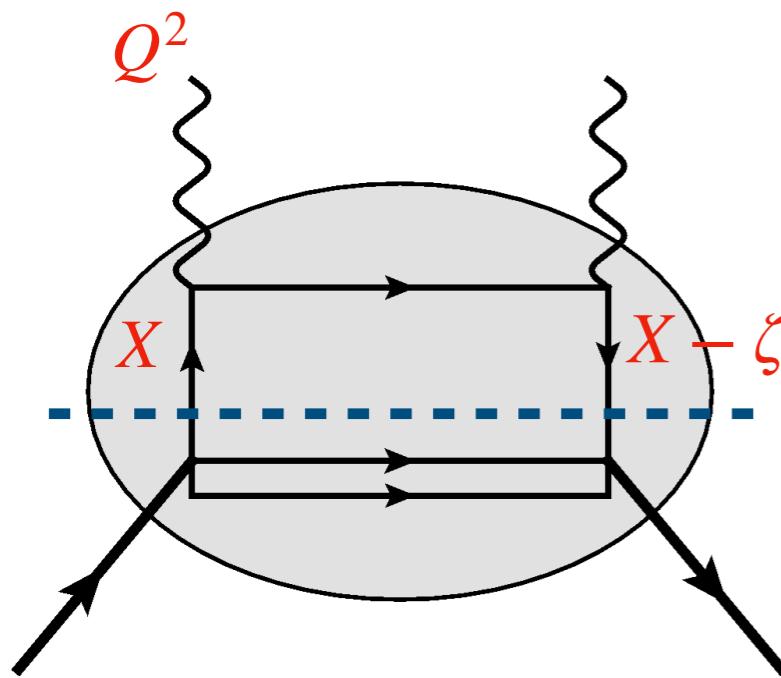
One Compton Form Factor



II. One-Loop Scalar Model

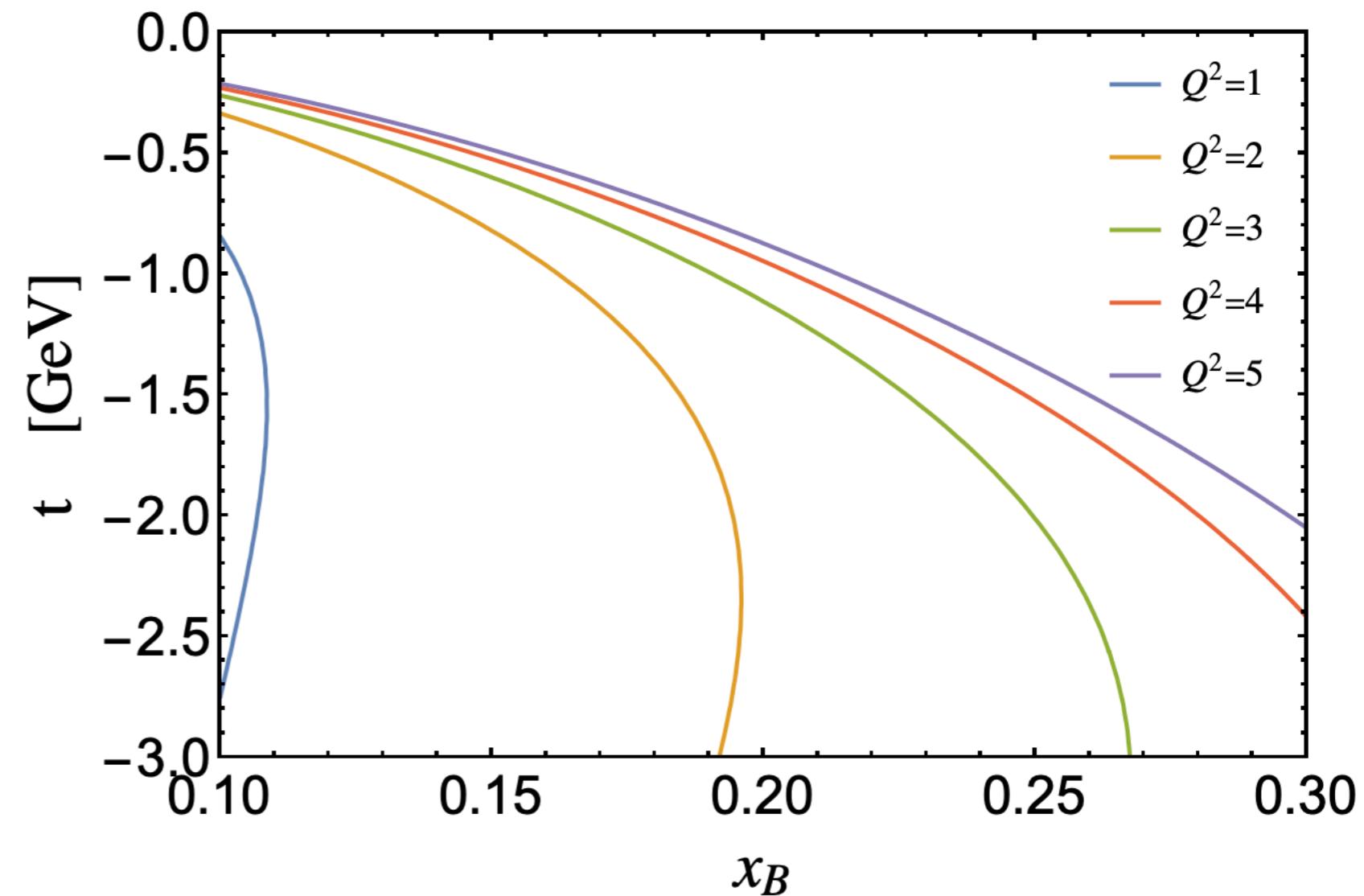
Kinematics

X : longitudinal momentum fraction,
 ζ : longitudinal momentum change of target,
 t : momentum transfer,



$$\text{CFF} \rightarrow \mathcal{F}(Q^2, x_B, t)$$

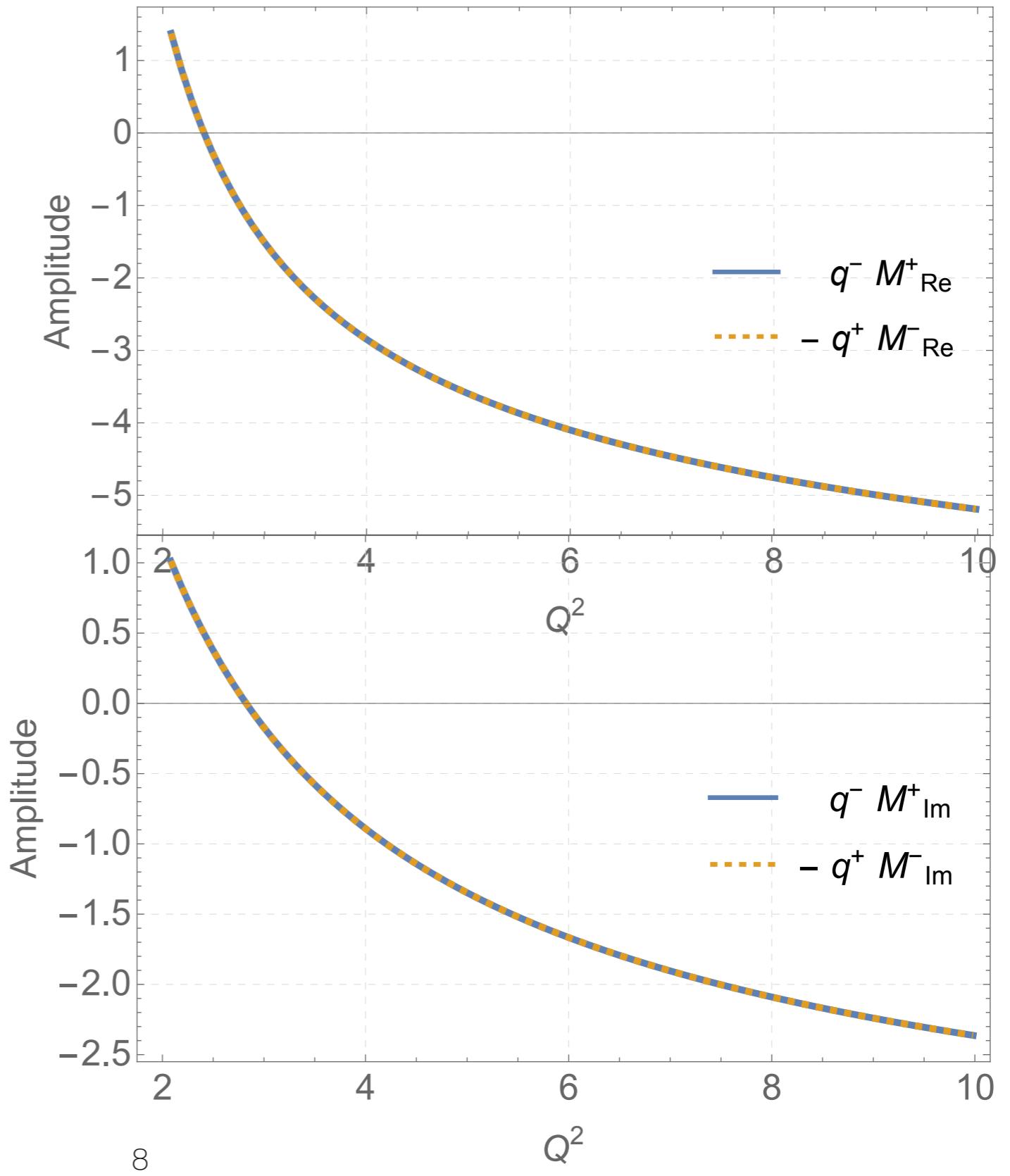
← GPD → $\mathcal{H}(X, \zeta, t)$ → $\lim_{Q^2 \rightarrow \infty} \zeta = x_B$
deeply virtual limit



Ward Identity

To check the numerical calculation,

$$q_\mu \mathcal{M}^\mu = q^+ \mathcal{M}^- + q^- \mathcal{M}^+ = 0$$



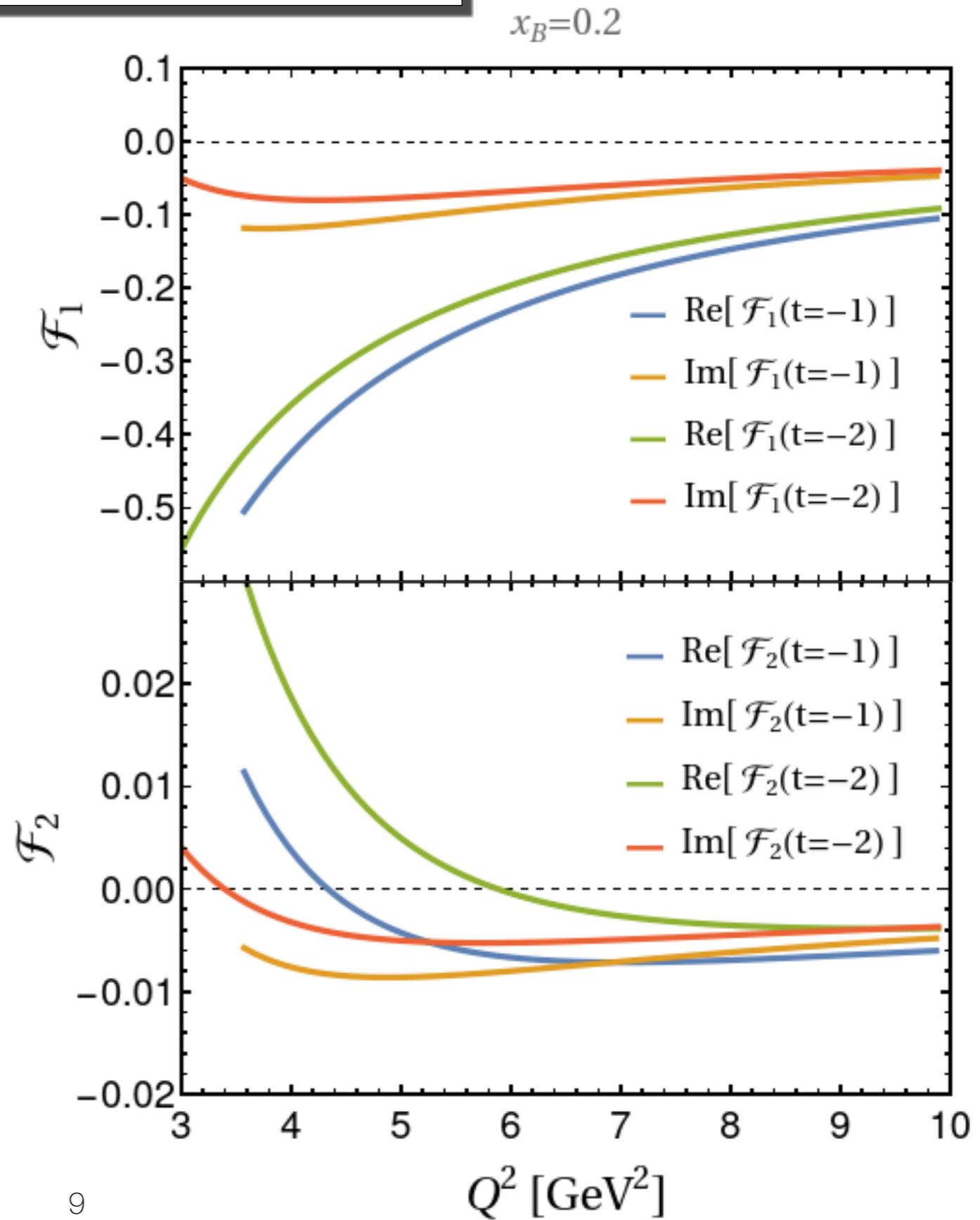
Compton Form Factors

Spin-0 Helium Target

$$\begin{aligned} \mathcal{M}^\mu = & \left[(\Delta \cdot q) q^\mu - q^2 \Delta^\mu \right] \mathcal{F}_1 \\ & + \left[(\Delta \cdot q) \mathcal{P}^\mu - (\mathcal{P} \cdot q) \Delta^\mu \right] \mathcal{F}_2 \end{aligned}$$

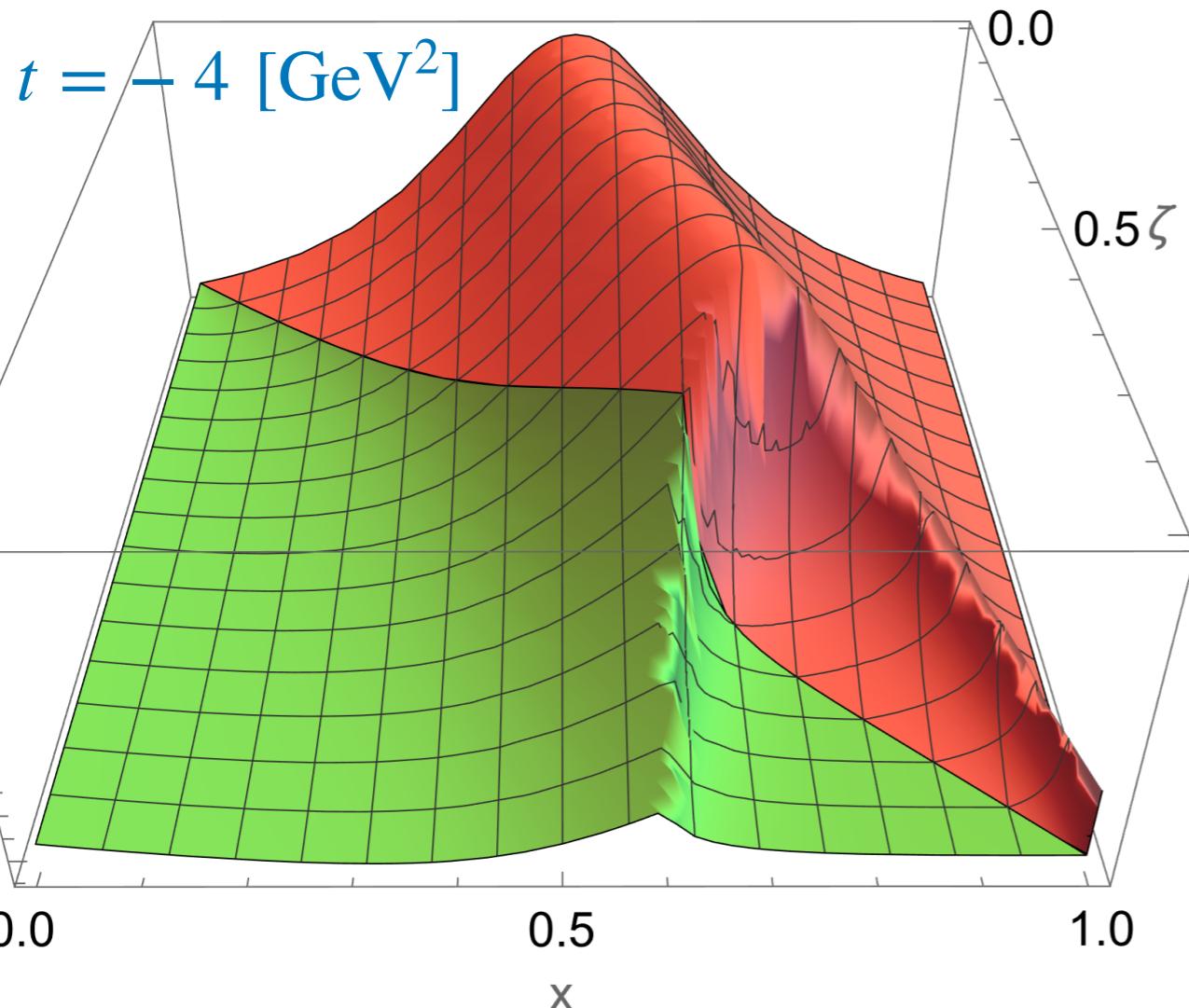
Remarks

1. \mathcal{F}_1 is approximately 10 orders of magnitude larger than \mathcal{F}_2 .
2. Size of CFFs seems to decrease As t gets larger.



III. Deeply Virtual Limit

Generalized Parton Distribution



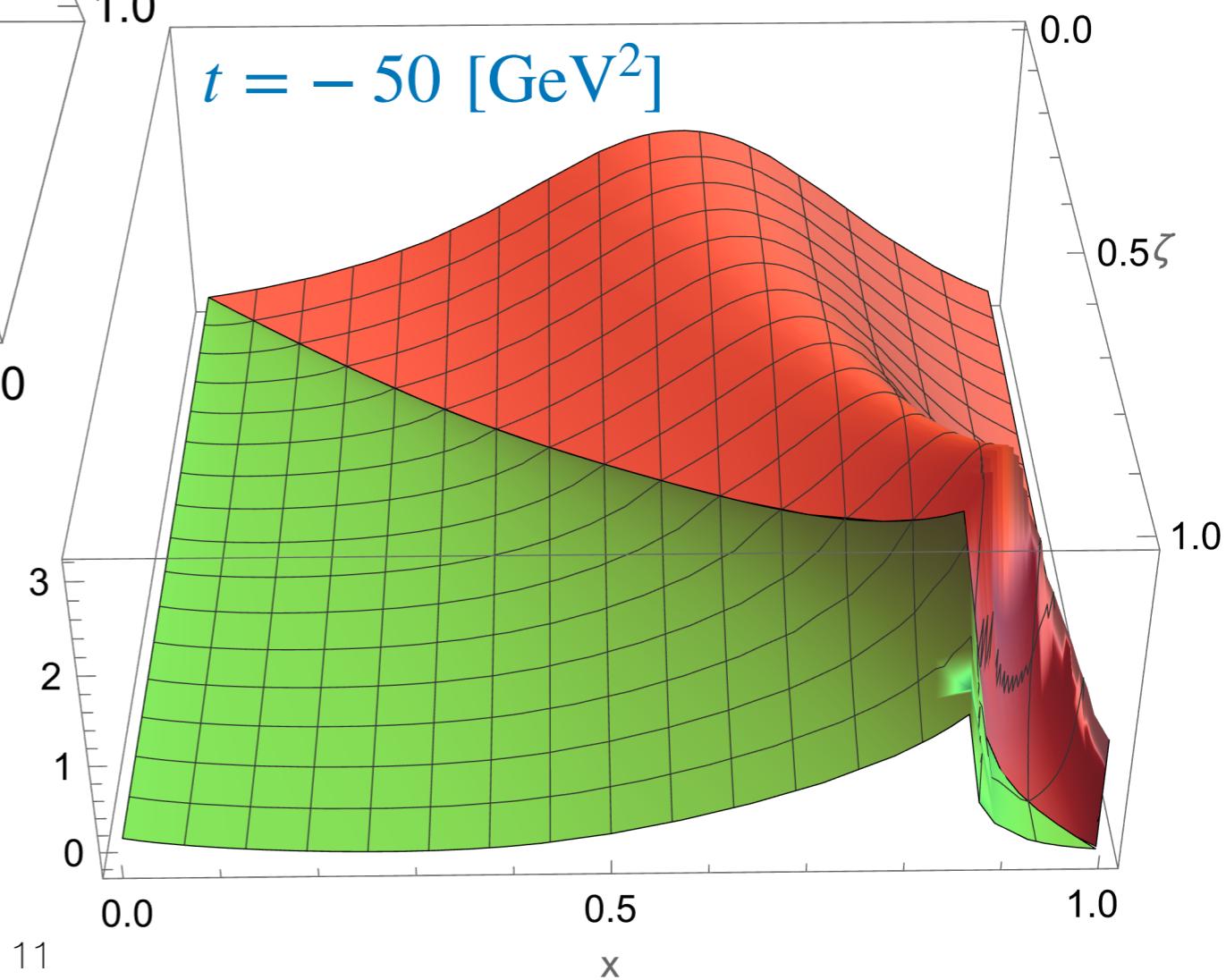
$$\mathcal{M}_{DVMP\ s+u}^+ \sim \int_0^1 \left(\frac{1}{x-\zeta} - \frac{1}{x} \right) H(x, \zeta, t)$$

$H_{\text{ERBL}} (0 \leq x \leq \zeta)$

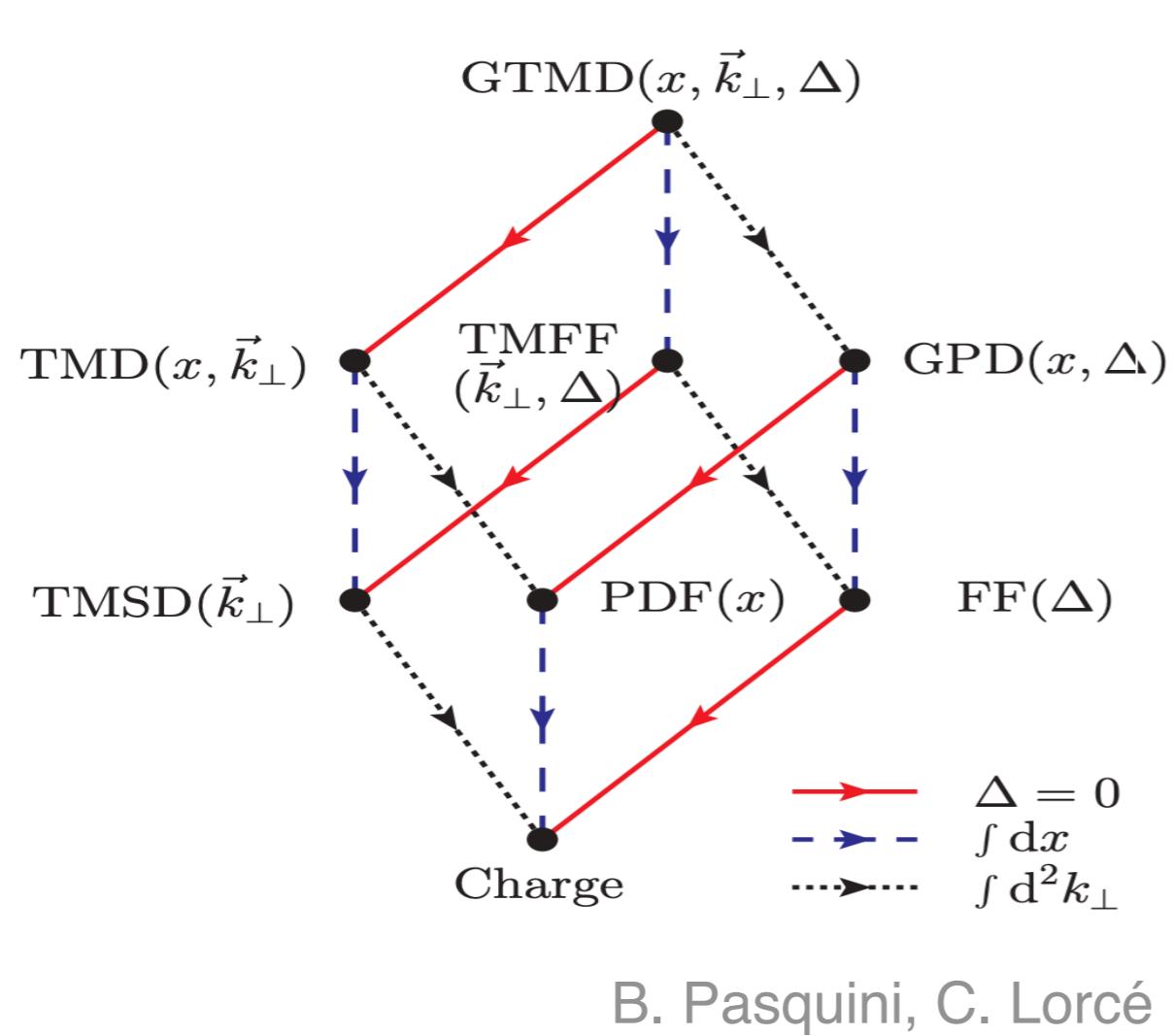
$H_{\text{DGLAP}} (\zeta \leq x \leq 1)$

Strange bump is probably
due to $(\zeta - 1)t - \zeta^2 M^2 < 0$.

$$|\Delta| = \sqrt{(\zeta - 1)t - \zeta^2 M^2}$$



GPD sum rule



$$\lim_{\xi, t \rightarrow 0} \int \frac{dy^-}{2\pi} e^{ixP^+y^-} \boxed{\langle P' | \bar{\psi}(0) \hat{\mathcal{O}} \psi(y) | P \rangle} \Big|_{y^+=y_\perp=0} = \int \frac{dy^-}{2\pi} e^{ixP^+y^-} \langle P | \bar{\psi}(0) \hat{\mathcal{O}} \psi(y) | P \rangle \Big|_{y^+=y_\perp=0}$$

$$\int dx \int \frac{dy^-}{2\pi} e^{ixP^+y^-} \boxed{\langle P' | \bar{\psi}(0) \hat{\mathcal{O}} \psi(y) | P \rangle} \Big|_{y^+=y_\perp=0} = \langle P' | \bar{\psi}(0) \hat{\mathcal{O}} \psi(0) | P \rangle \Big|_{y^+=y_\perp=0}$$

Parton Distribution Functions :

$$H^q(x, 0, 0) = f_1(x), \quad \tilde{H}^q(x, 0, 0) = g_1(x),$$

Form Factors (first Mellin moment) :

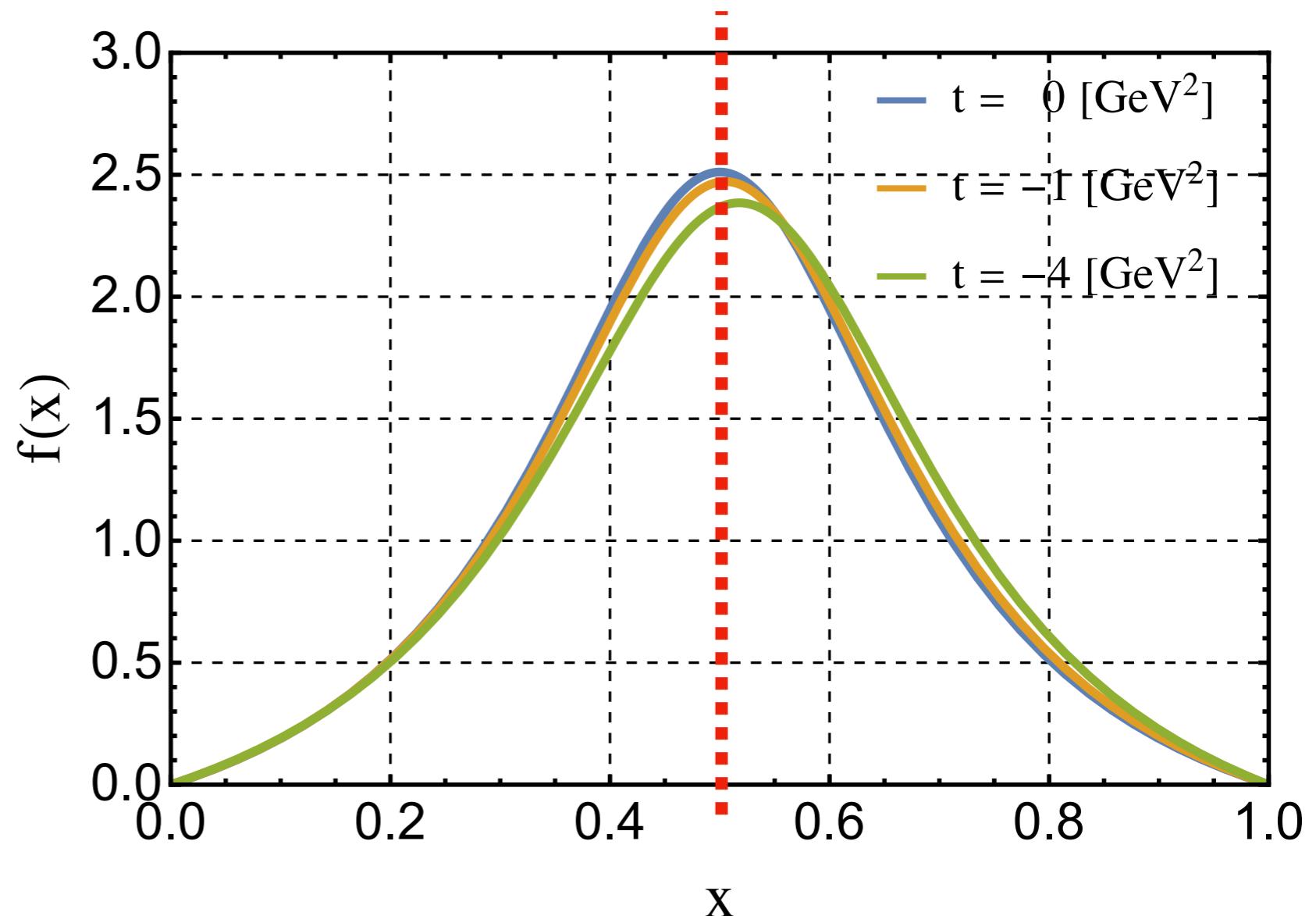
$$\int dx H^q(x, \xi, t) = F_1^q(t), \quad \int dx E^q(x, \xi, t) = F_2^q(t)$$

$$\int dx \tilde{H}^q(x, \xi, t) = G_A^q(t), \quad \int dx \tilde{E}^q(x, \xi, t) = G_P^q(t)$$

Parton Distribution Function

$$\lim_{\xi,t \rightarrow 0} H(x, \xi, t) = \lim_{\zeta,t \rightarrow 0} H(X, \zeta, t) = H(X, 0, 0) = f(X)$$

1. $\zeta \rightarrow 0$ corresponds to
 $t \rightarrow 0$ in (1+1). In (3+1),
PDF can be obtained with
 $\zeta, t \rightarrow 0$ simultaneously.
2. Helium is consist of **two effective quarks** with
equal masses in our model.
Momentum fraction of a
single quark has highest
probability at 0.5.



Electromagnetic Form Factor

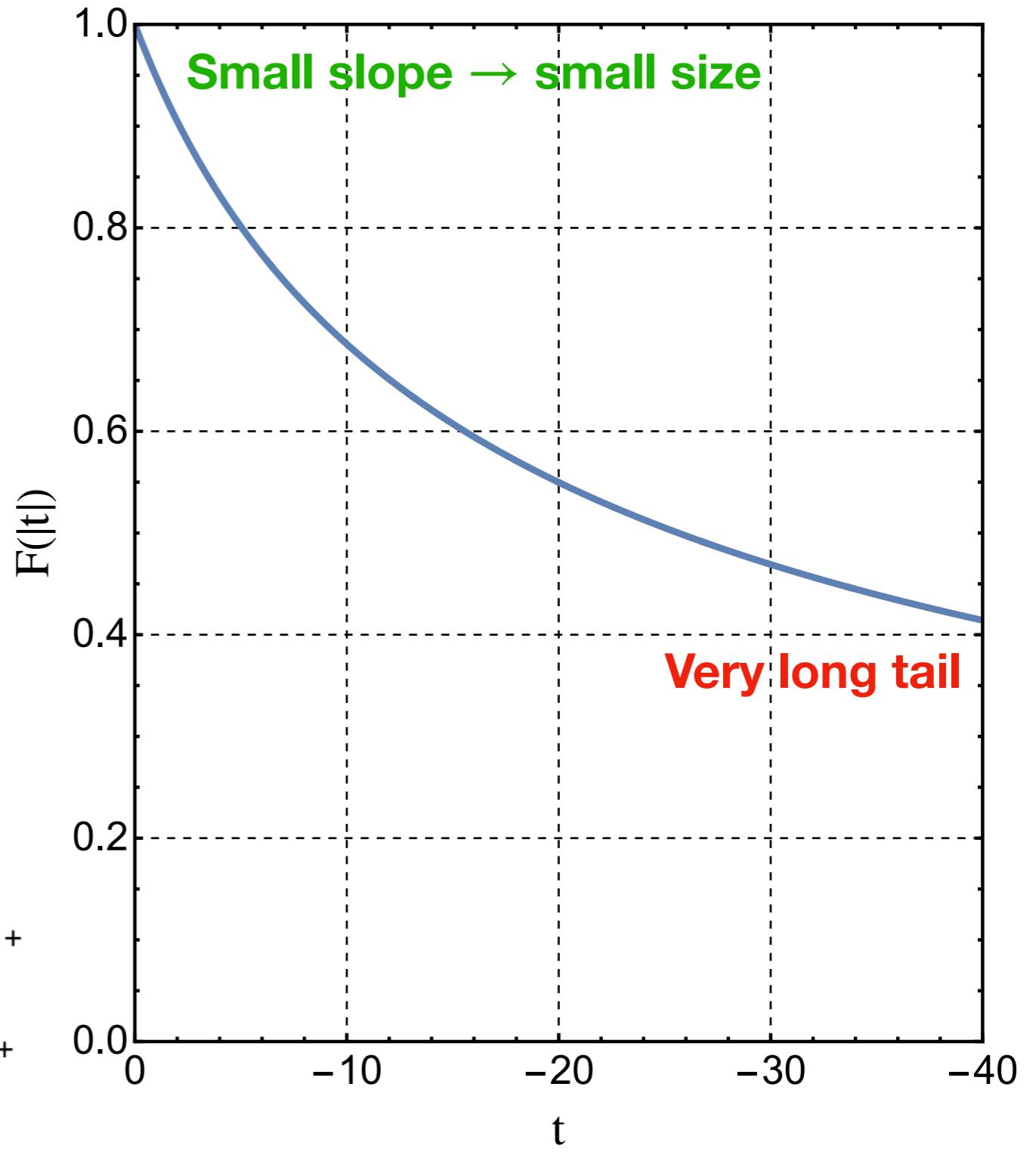
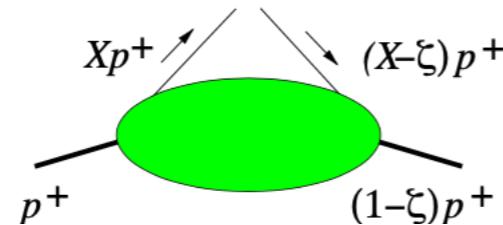
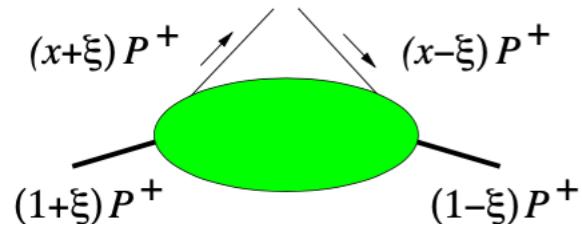
1st Melin moment

$$\int dx H(x, \xi, t)$$

$$= \int dX \frac{2}{2 - \zeta} H(X, \zeta, t) = F(t),$$

with different conventions,

$$X = \frac{x + \xi}{1 + \xi}, \quad \zeta = \frac{2\xi}{1 + \xi}$$



Energy Momentum Tensor FF

2nd Melin moment

$$\int dx \ xH(x, \xi, t) = A(t) + 4\xi^2 C(t)$$

$$= \int dX \frac{2(2X - \zeta)}{(2 - \zeta)^2} H(X, \zeta, t)$$

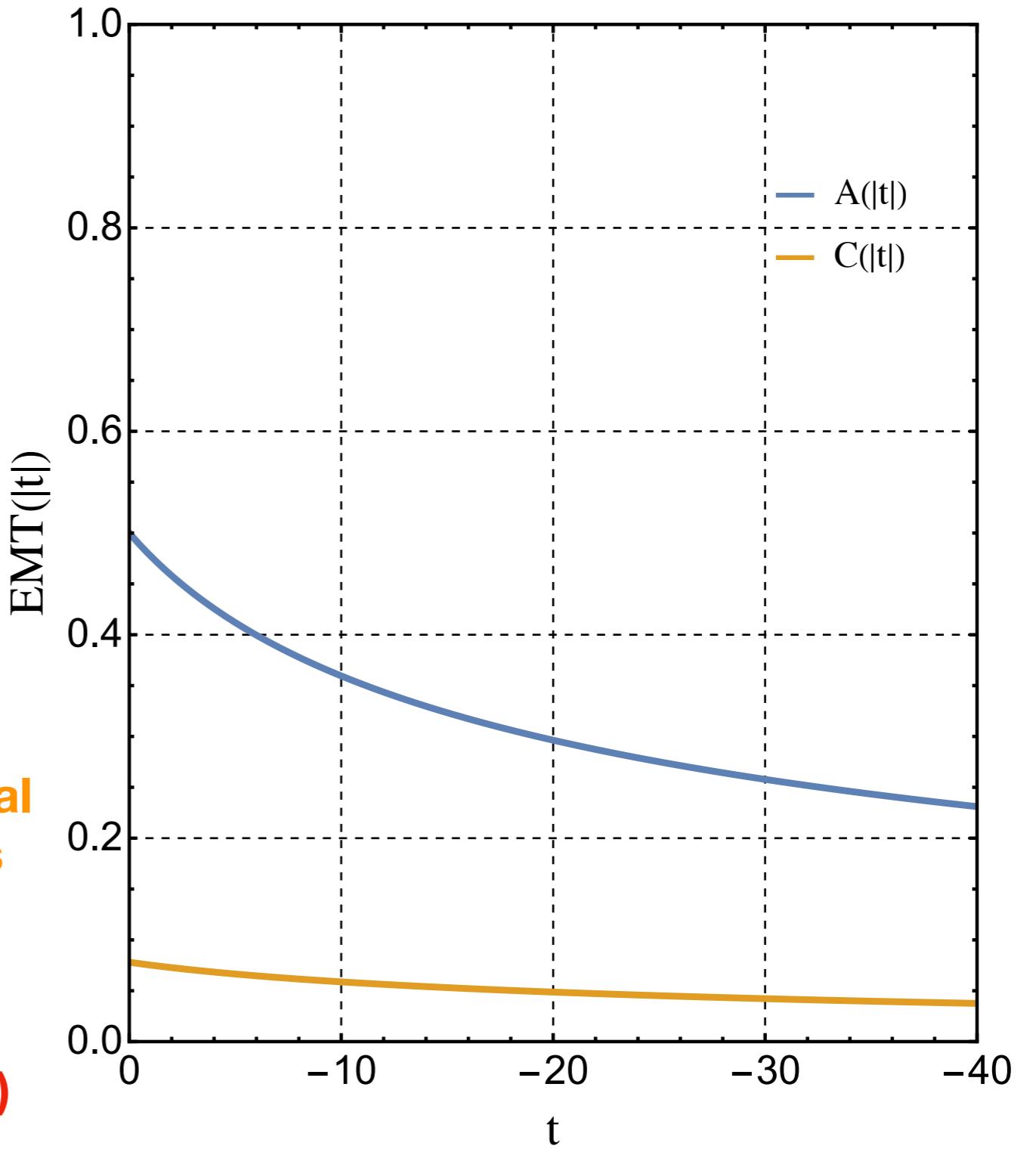
$$= A(t) + 4 \left(\frac{\zeta}{2 - \zeta} \right)^2 C(t)$$

↓
Mass

(half mass of quark)

↓
Mechanical properties

(pressure & shear forces)



IV. One-loop Scalar Model vs Deeply Virtual Limit

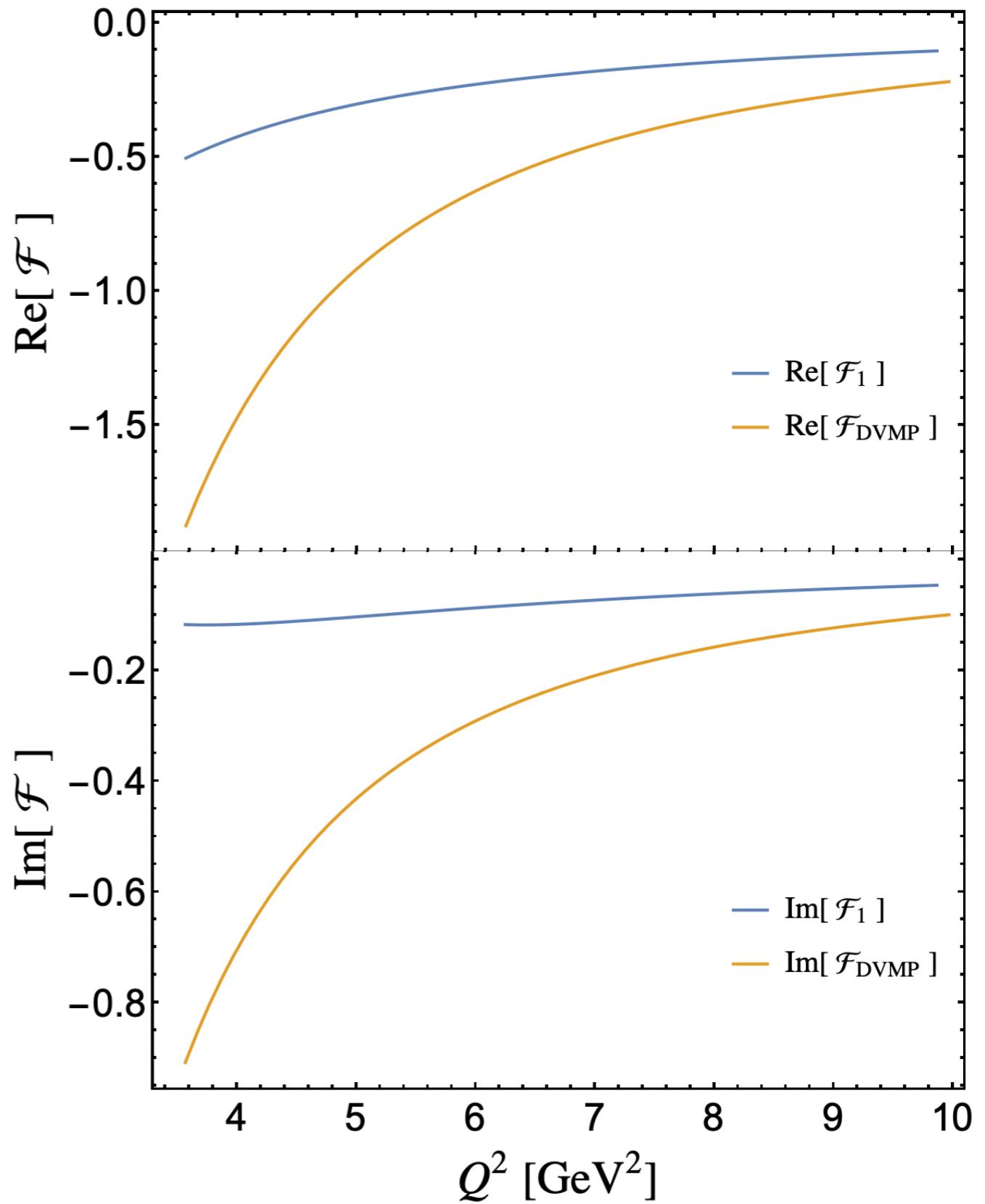
Comparison two models for CFFs

One-loop scalar model :

$$\mathcal{F}_1 \gg \mathcal{F}_2$$

Deeply virtual limit :

only \mathcal{F}_{DVMP}



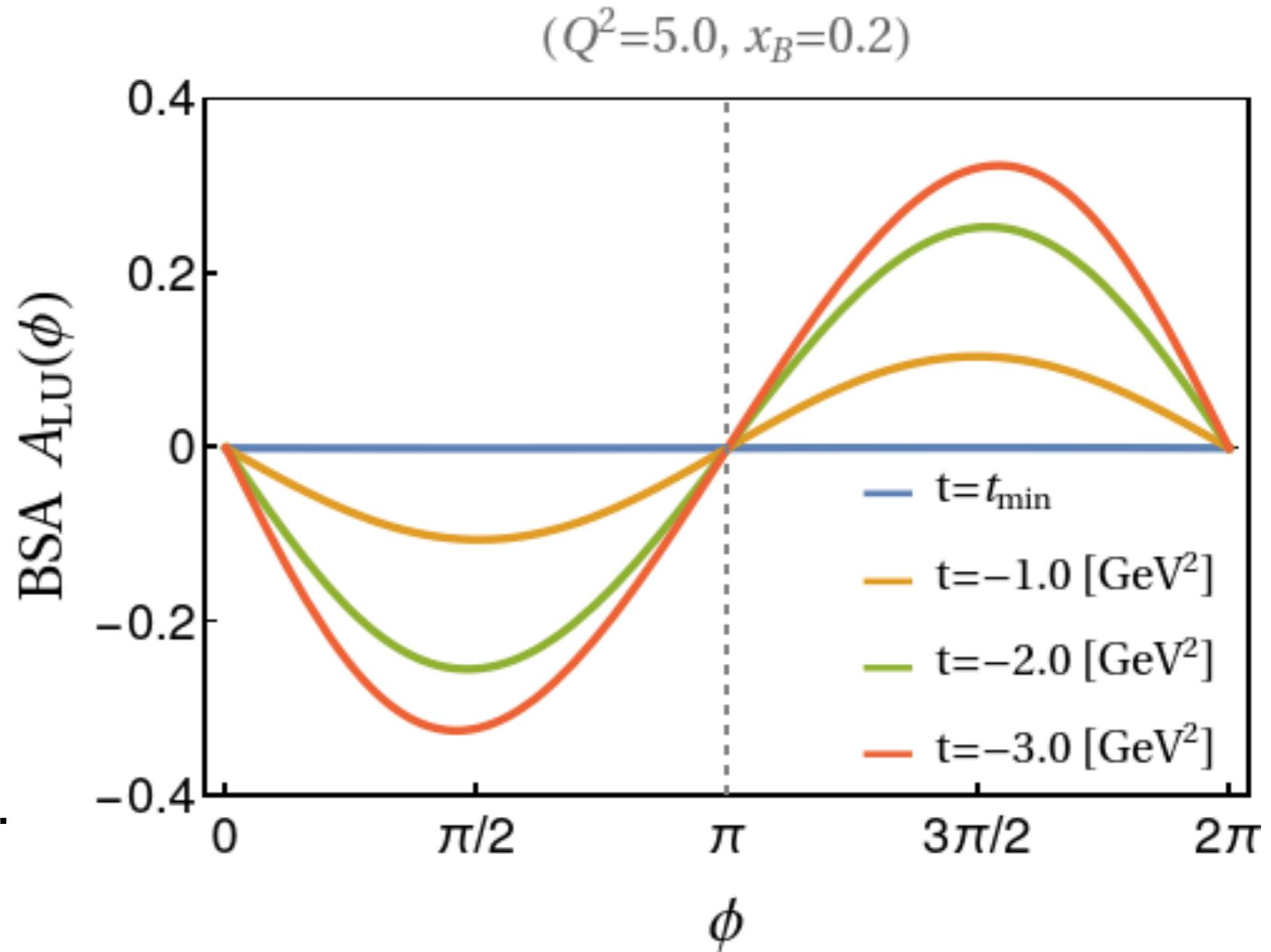
Beam Spin Asymmetry

$$A_{LU}^S(\phi) = \frac{d\sigma_{BSA}^S}{d\sigma_T^S (1 + \epsilon \cos(2\phi)) + d\sigma_L^S \epsilon_L + d\sigma_{LT}^S \cos(\phi)\sqrt{\epsilon_L(1 + \epsilon)/2}}$$

$$d\sigma_{BSA}^S = S_A (\mathcal{F}_1 \mathcal{F}_2^* - \mathcal{F}_1^* \mathcal{F}_2)$$

Remarks

1. As t increases, magnitude of BSA increases.
2. For a given Q^2 and x_B , if $t = t_{min}$, BSA vanishes.





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"Thank you for listening."