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# Algebraic Methods for Efficient Hamiltonian Simulation via Quantum Computers

Efekan Kökcü

Lawrence Berkeley National Laboratory

[ekokcu@lbl.gov](mailto:ekokcu@lbl.gov)



# Outline

- Quantum computing
  - Qubits and quantum gates
- Quantum simulation of spin-1/2 systems via quantum computers
  - Simulation overall: state prep, time evolution, measurement
  - Time evolution: Trotter-Suzuki approach
- Beyond spin 1/2 systems:
  - Fermions
  - Bosons
- Algebraic Compression of Free Fermionic Evolution

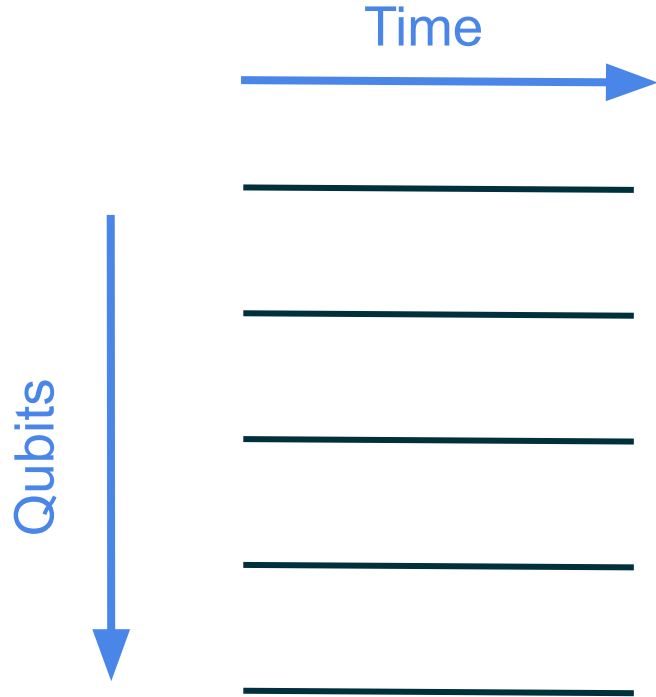


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# Quantum Computing: Qubits

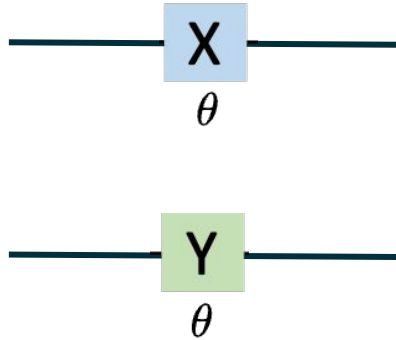


- Qubits are two level systems which are labeled as 0 and 1
- $n$  qubits can store  $2^n$  complex numbers at once! Thus, quantum computers are powerful candidates for quantum simulation



# Quantum Computing: Quantum gates

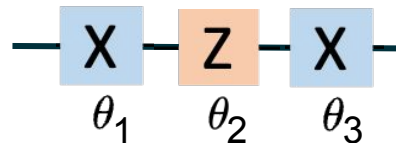
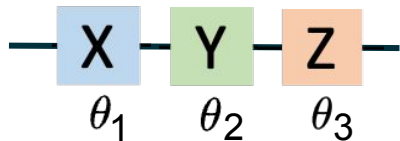
- Commonly, 1-qubit gates are defined as exponentials of Pauli matrices



# Quantum Computing: Quantum gates

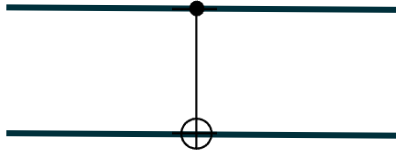
- Commonly, 1-qubit gates are defined as exponentials of Pauli matrices

- Any 1-qubit unitary can be written as follows:



# Quantum Computing: Quantum gates

- Controlled not gate is defined as follows:

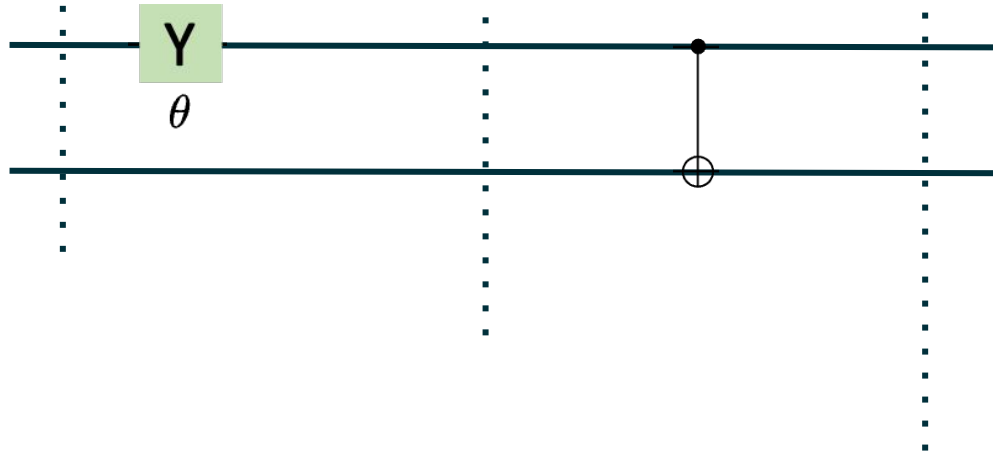


- It flips the target qubit if the control qubit is in 1 state
- Currently, these are more expensive and more noisy. In the fault tolerant era, the 1-q gates are going to be more expensive.



# Quantum Computing: Quantum gates

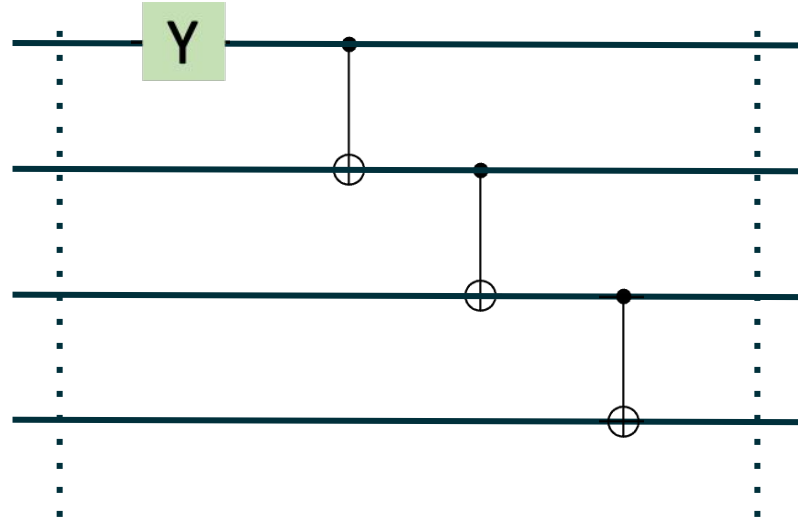
- Example: the following generates Bell state for





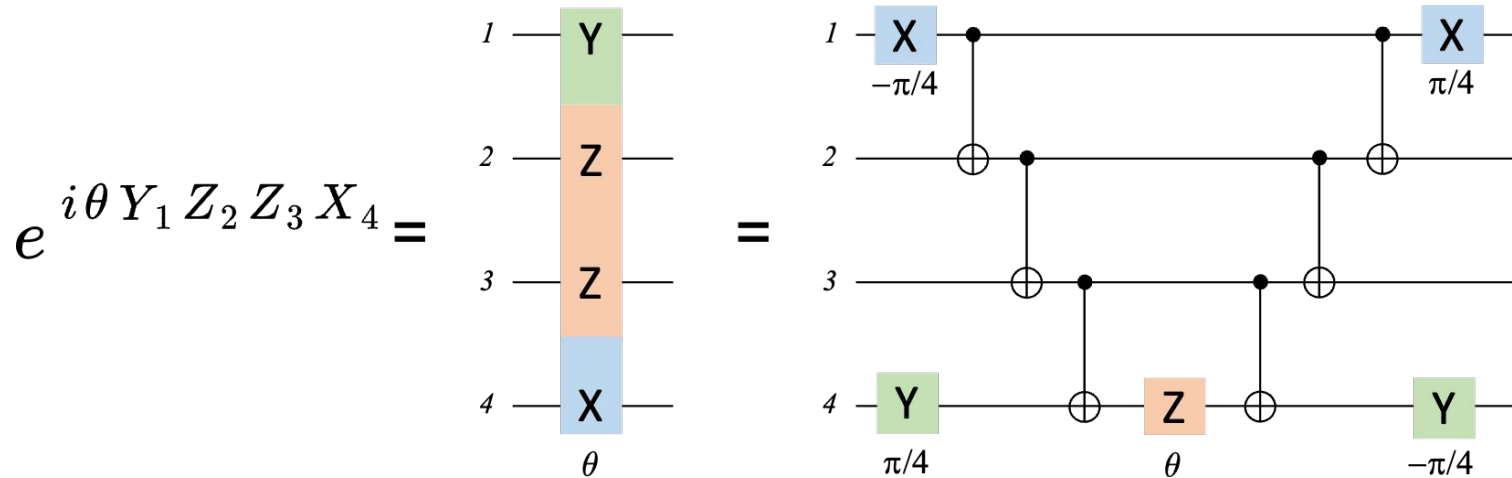
# Quantum Computing: Quantum gates

- Similarly, the GHZ state can be generated as follows:



# Quantum Computing: Quantum gates

- Using these gates, we can implement rotations of any tensor product of Pauli matrices



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# Quantum simulation of spin 1/2 systems

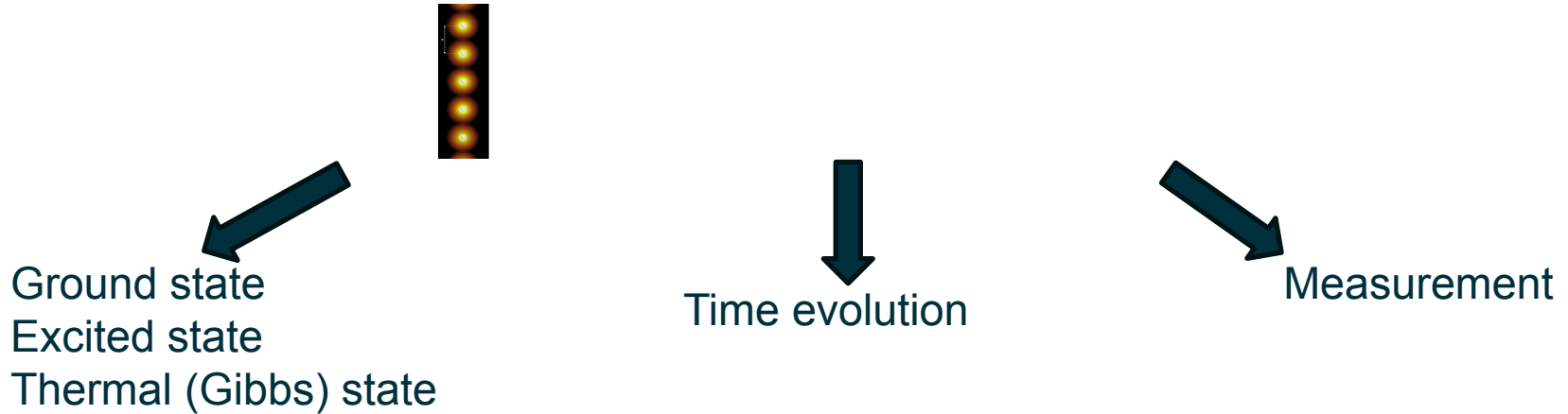
- To simulate physical models, one need to map physical degrees of freedom to the d.o.f. of the quantum computer
- For spin 1/2 systems, we can map as follows:

$$|0\rangle \equiv |\uparrow\rangle \quad |1\rangle \equiv |\downarrow\rangle$$

- Then, each spin operator correspond to the same 1-q quantum gate (to a factor of 2)!



# Quantum simulation via quantum computing



# Time evolution: Trotter-Suzuki decomposition

- Consider Kitaev spin chain on 5 spins:

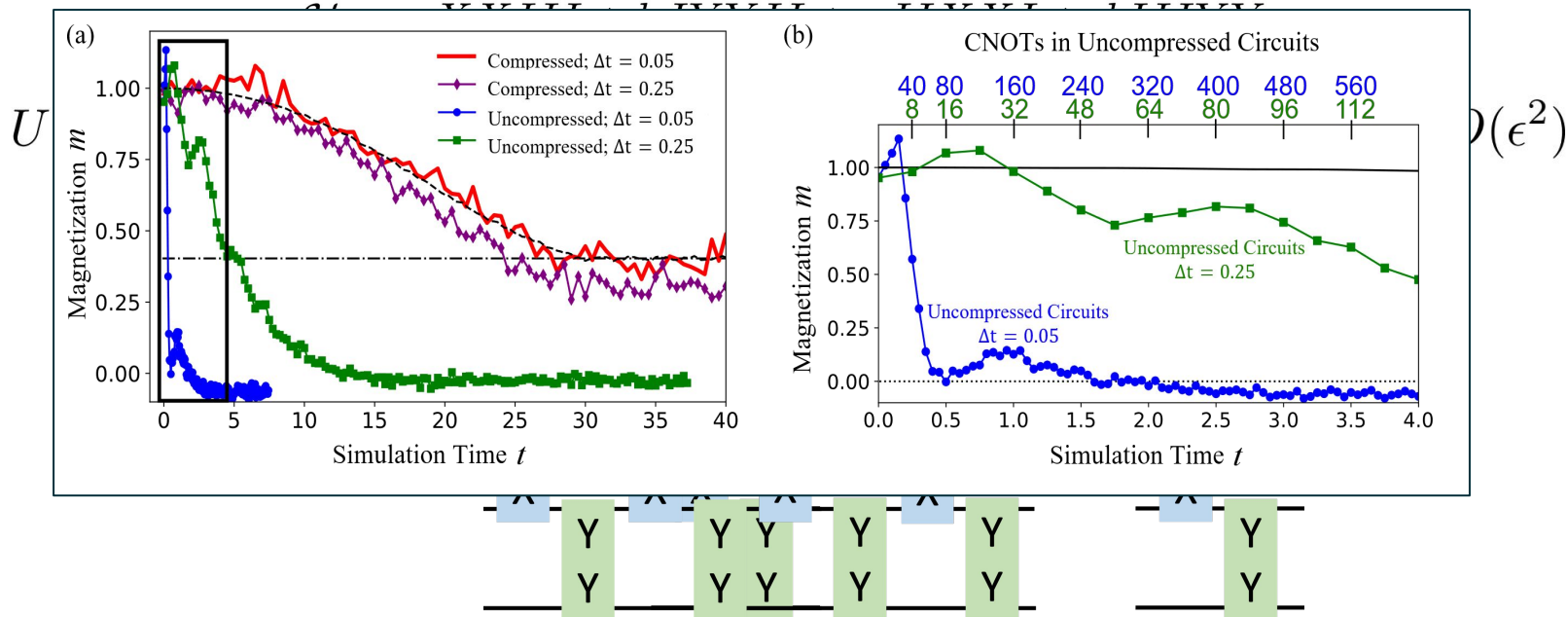
$$\mathcal{H} = a XXIII + b IYYII + c IIXXI + d IIIYY$$

$$U(t) = e^{-it\mathcal{H}} \neq e^{-ita XXIII} e^{-itb IYYII} e^{-itc IIXXI} e^{-itd IIIYY}$$



# Time evolution: Trotter-Suzuki decomposition

- Consider Kitaev spin chain on 5 spins:



# Outline

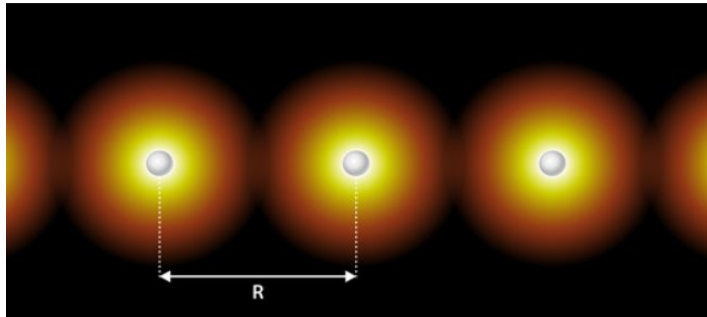
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# Beyond spin 1/2 systems: Fermions

- Fermionic systems in condensed matter consist of the following states



$|\text{vac}\rangle$

$|\text{vac}\rangle$

$|\text{vac}\rangle$

$c^\dagger|\text{vac}\rangle$

$c^\dagger|\text{vac}\rangle$

$c^\dagger|\text{vac}\rangle$

- These states can directly be mapped to qubits:

$$|0\rangle \equiv |\text{vac}\rangle$$

$$|1\rangle \equiv c^\dagger|\text{vac}\rangle$$

- One needs to make sure that the fermions generate a minus sign under exchange.



# Beyond spin 1/2 systems: Fermions

- One example of these mappings is the Jordan-Wigner mapping:
- There are other mappings such as Bravyi-Kitaev, and generic ternary tree mappings [1,2]

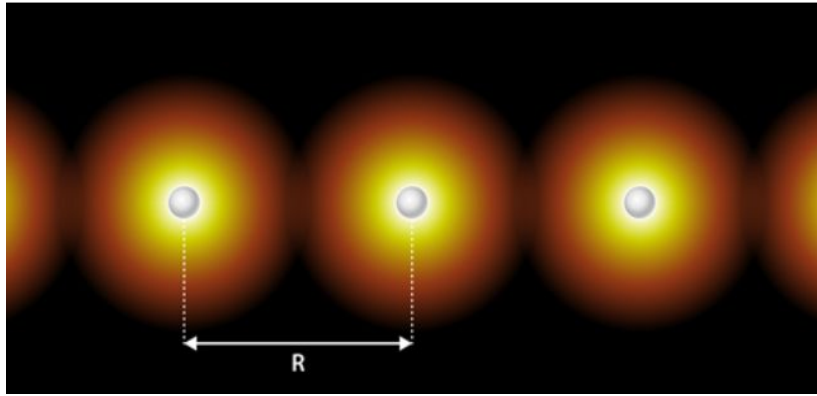
[1] A. Miller et al (2023), PRX Quantum

[2] Y. Liu et al (2025), IEEE Int. Symp. HPCA



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# Beyond spin 1/2 systems: Bosons



$|\text{vac}\rangle$

$|\text{vac}\rangle$

$|\text{vac}\rangle$

⋮

⋮

⋮

[1] R. Somma et al. Inter. Jour. Quant. Inf. 1.02 (2003): 189-206

[2] A. Miessen et al, PRR 3.4 (2021): 043212

- Boson modes can be seen as quantum harmonic oscillators.
- If truncated in  $N$  states, each boson mode can be mapped to  $N$  qubits (unary mapping) or  $\log_2 N$  qubits (binary mapping)
- Because the qubit operators commute, bosonic statistics is readily satisfied!

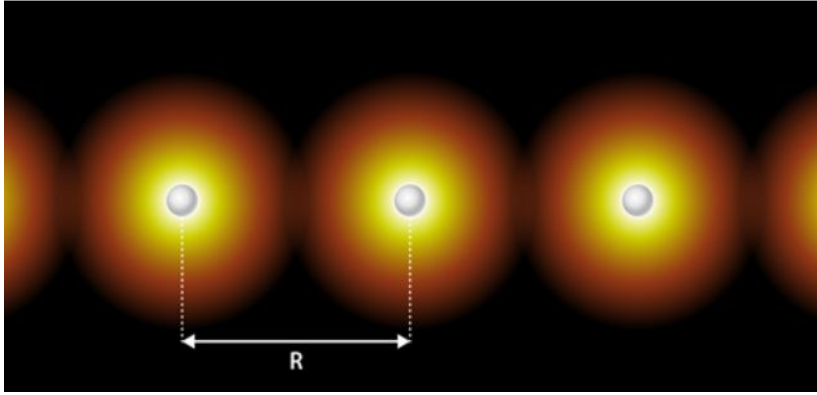
[3] N. Sawaya et al, npjQI 6.1 (2020): 49

[4] B. Peng et al, Quant. Sci. Tech. 10.2 (2025):023002



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# Beyond spin 1/2 systems: Bosons



- Boson modes can be seen as quantum harmonic oscillators.
- Jordan-Lee-Preskill mapping maps each state to an n-qubit state, digitizing the field space.

[1] N. Klyco et al, PRA 99.5 (2019): 052335

[2] R. Ferrel et al, PRD 109.11 (2024): 114510

[3] S. P. Jordan, K. S. M. Lee, and J. Preskill (2011), [Quant. Inf. Comput.14,1014(2014)], 1112.4833



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# Acknowledgements



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E. Kökcü et al (2022), PRA 105(3), 032420

D. Camps et al (2022), SIAM, 43(3), 1084-1108.

E. Kökcü et al (2023), arXiv:2303.09538



Wibe de Jong  
LBNL



Alexander Kemper  
NCSU

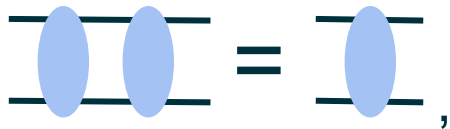


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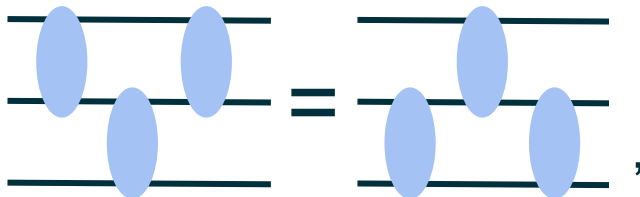
# Algebraic Compression

Consider quantum gates that satisfy the following relations:

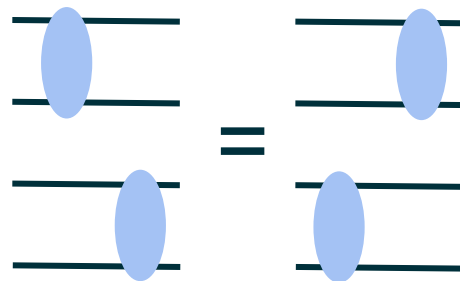
Fusion



Turnover



Commutation



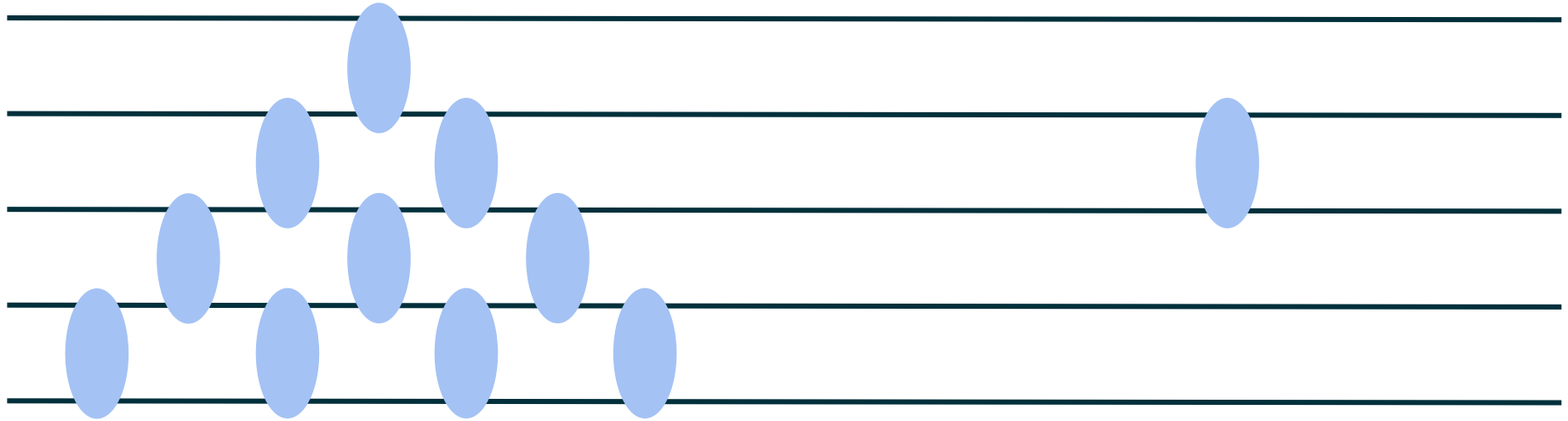
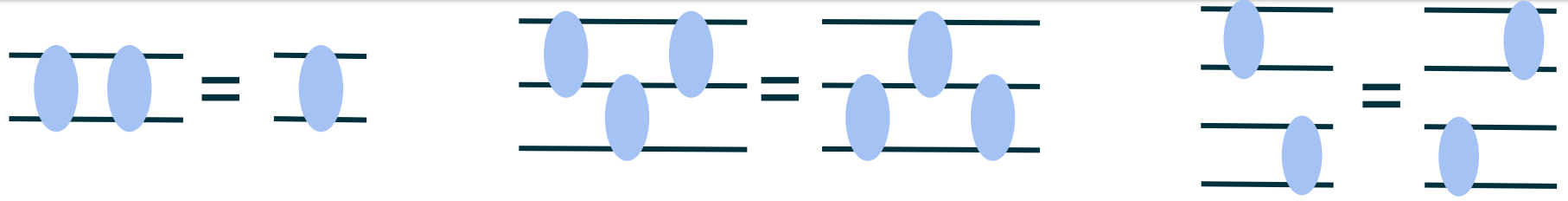
We call them **blocks**. Sequences of blocks can be simplified via algebraic compression [1-2]

[1] E. Kökcü et al (2022), PRA 105(3), 032420

[2] D. Camps et al (2022), SIAM, 43(3), 1084-1108.

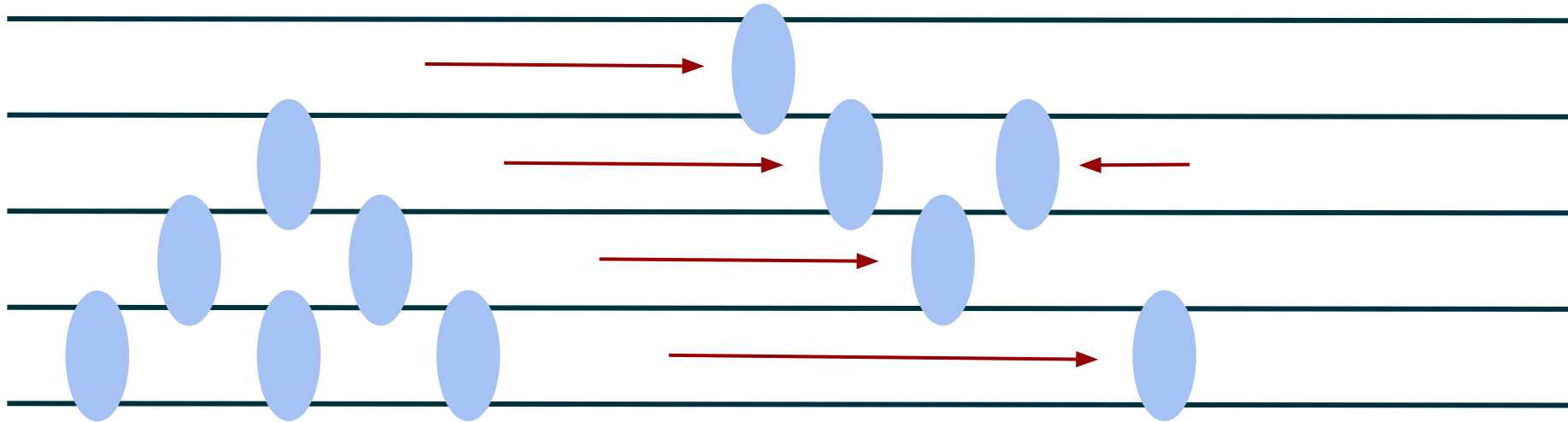
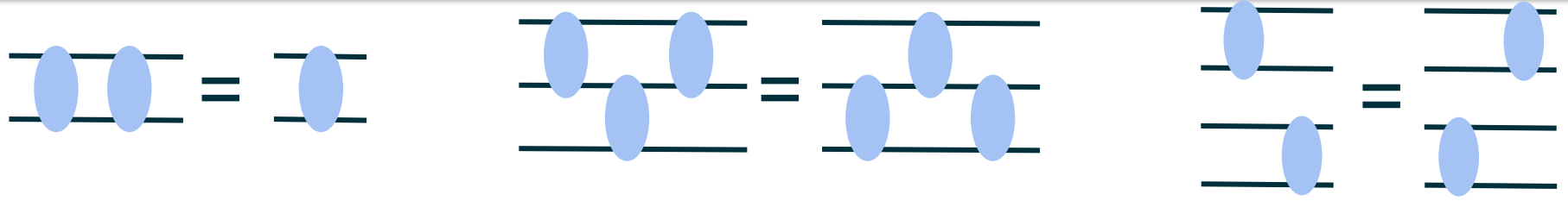


# Simplification of blocks

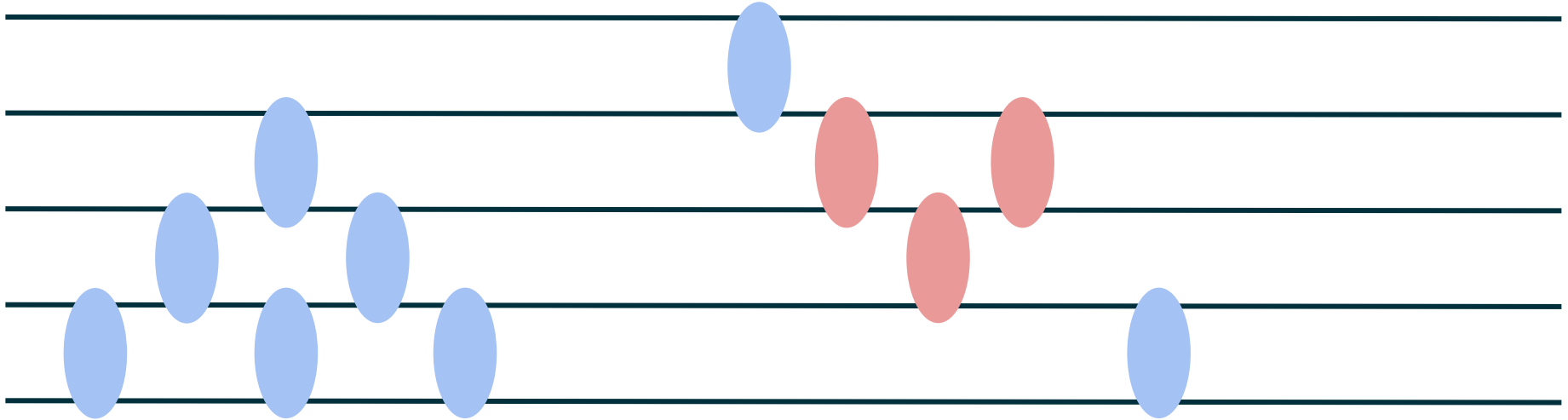
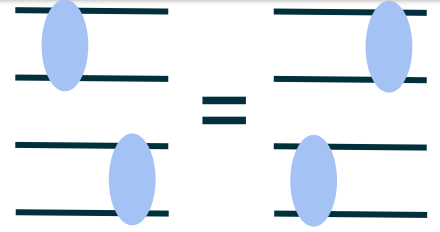
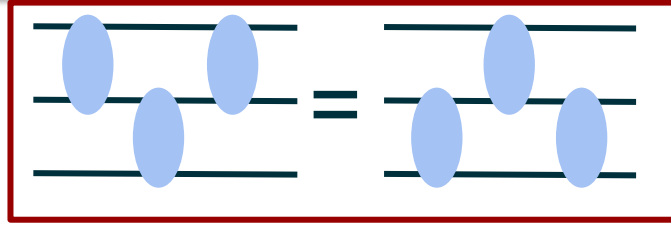




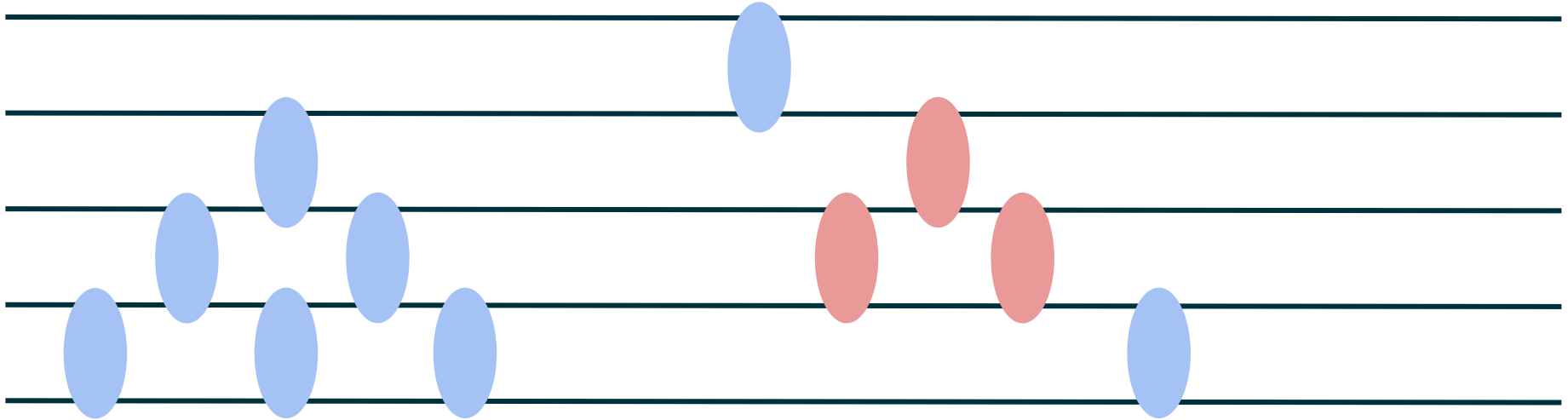
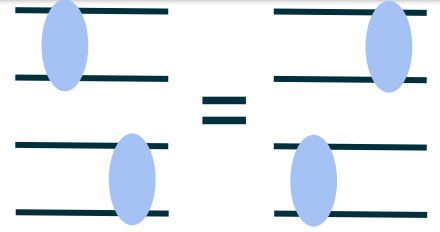
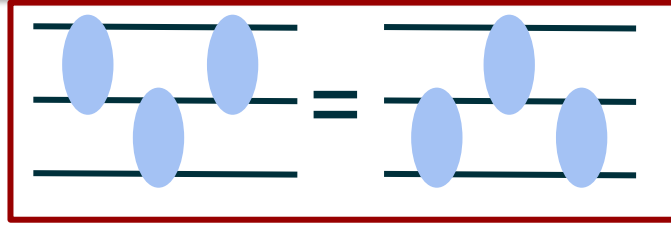
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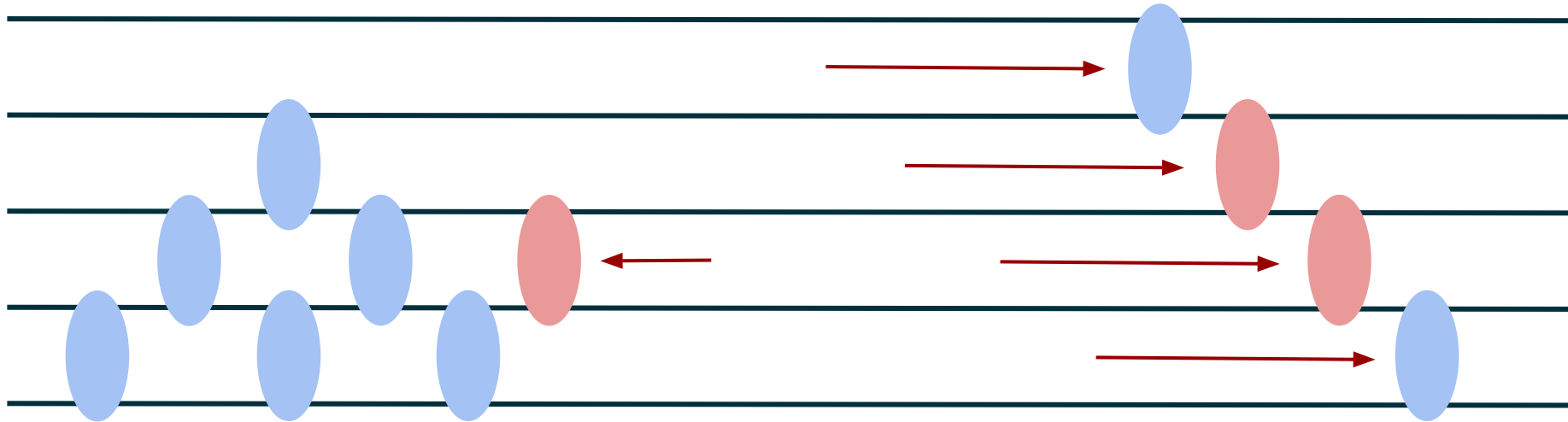
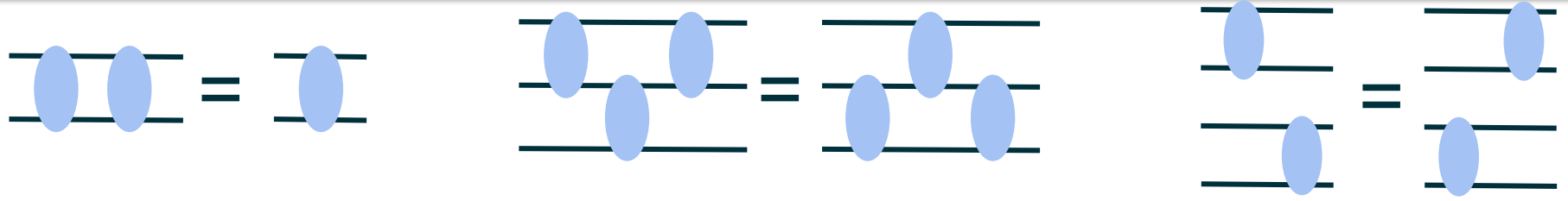
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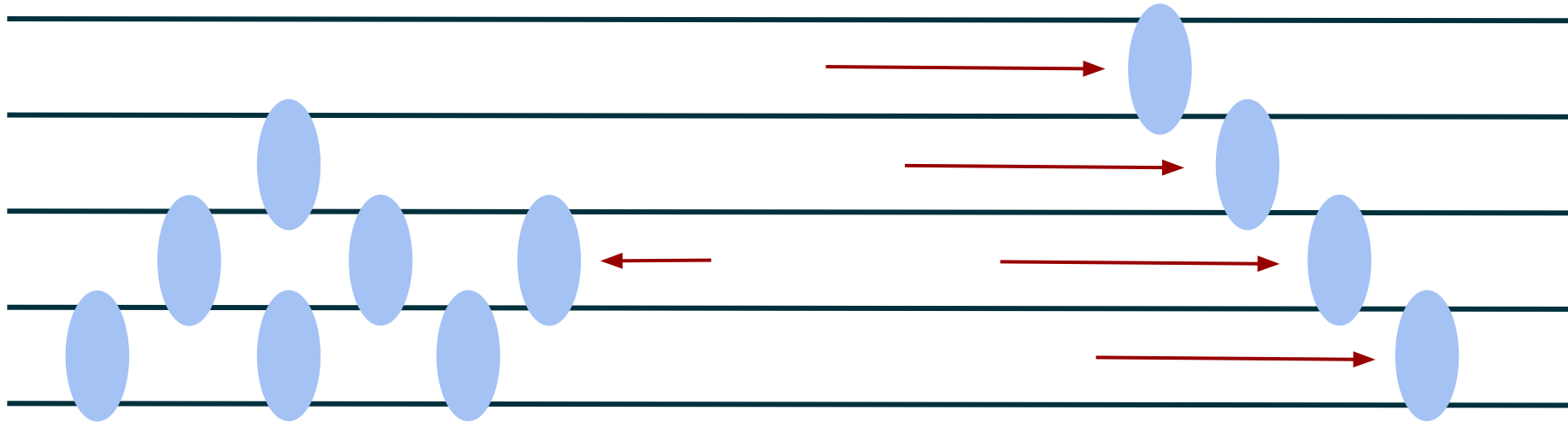
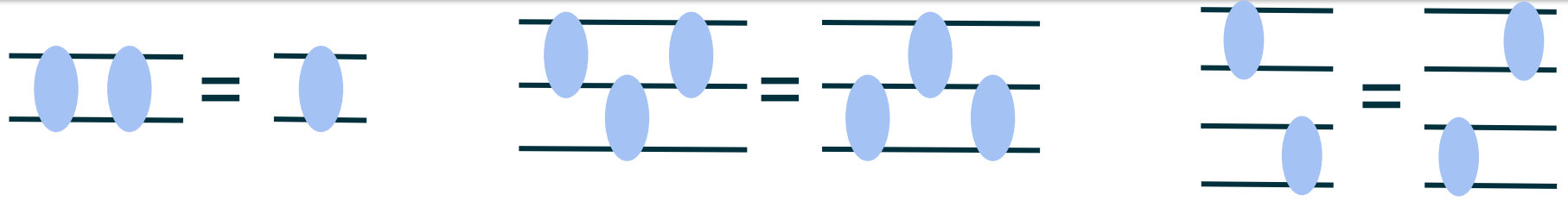
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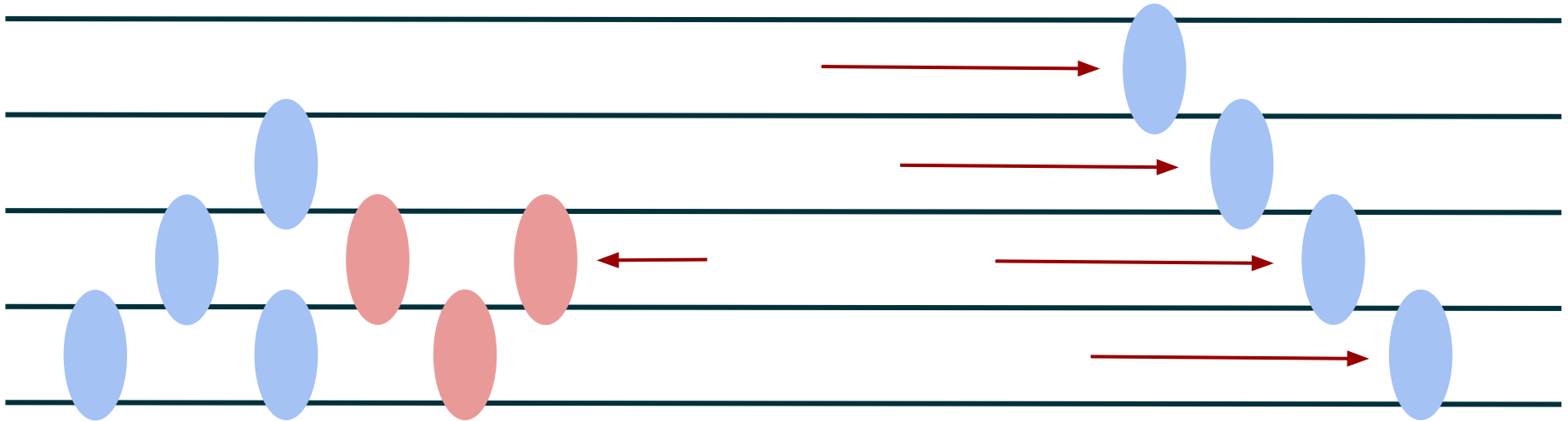
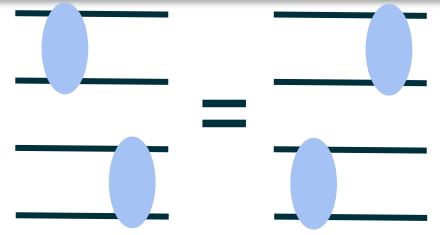
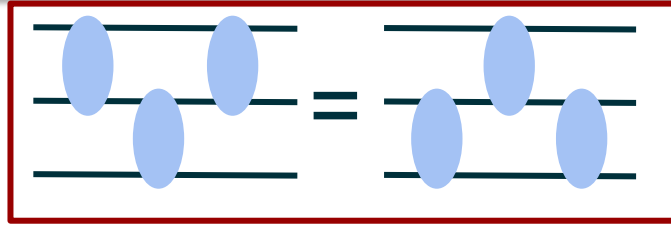
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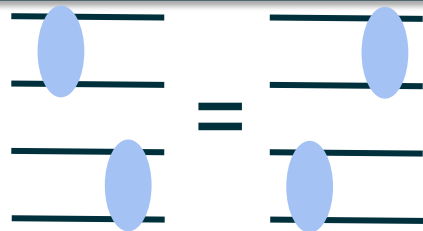
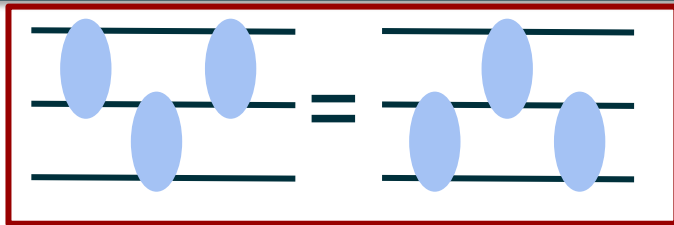
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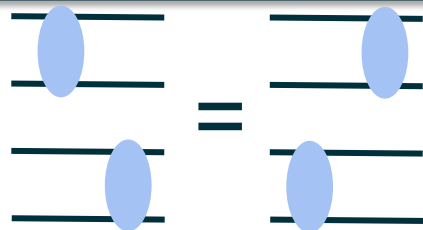
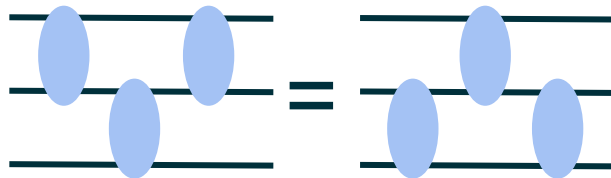
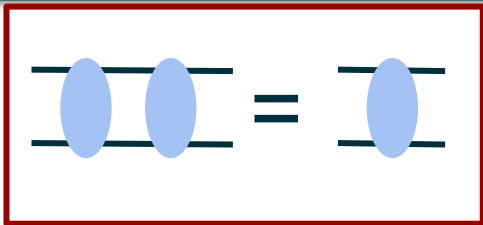
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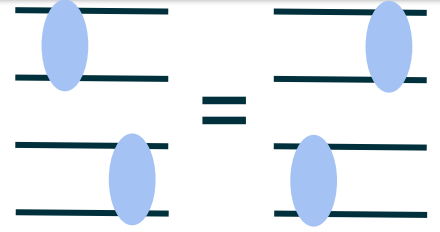
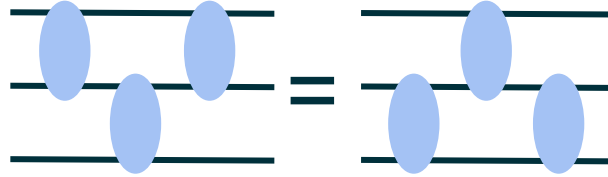
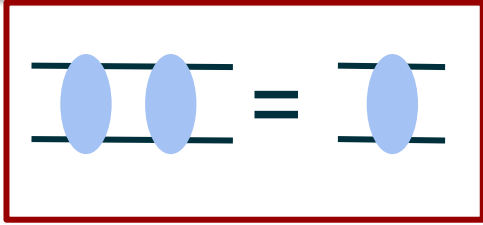


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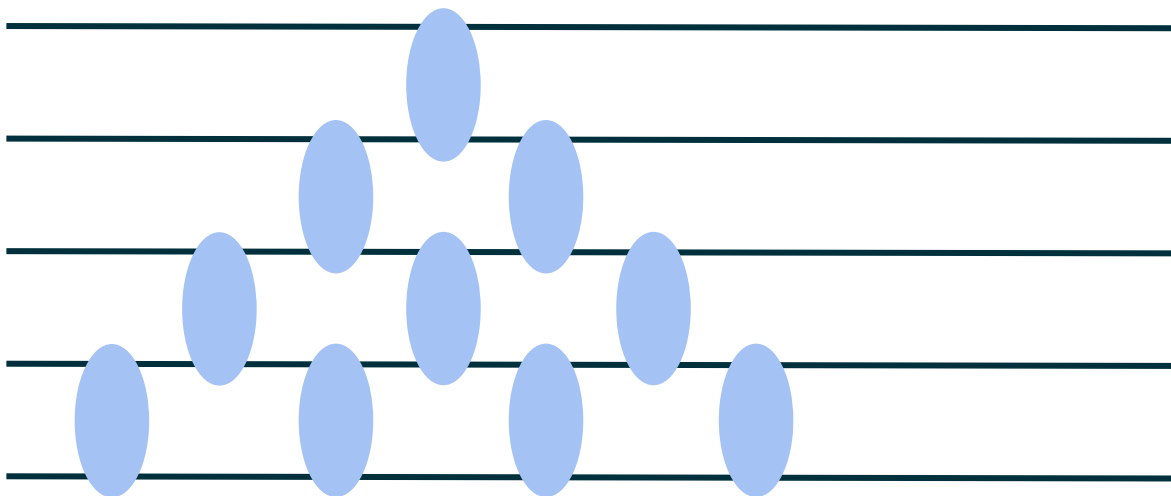
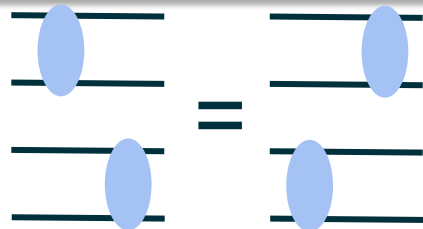
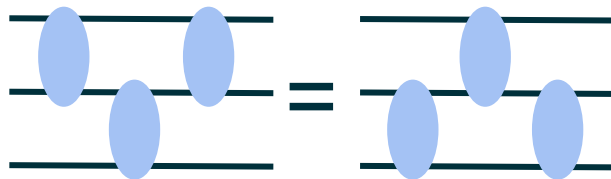
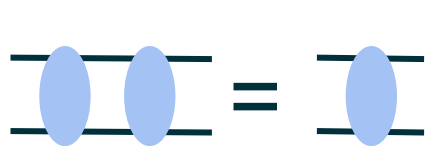




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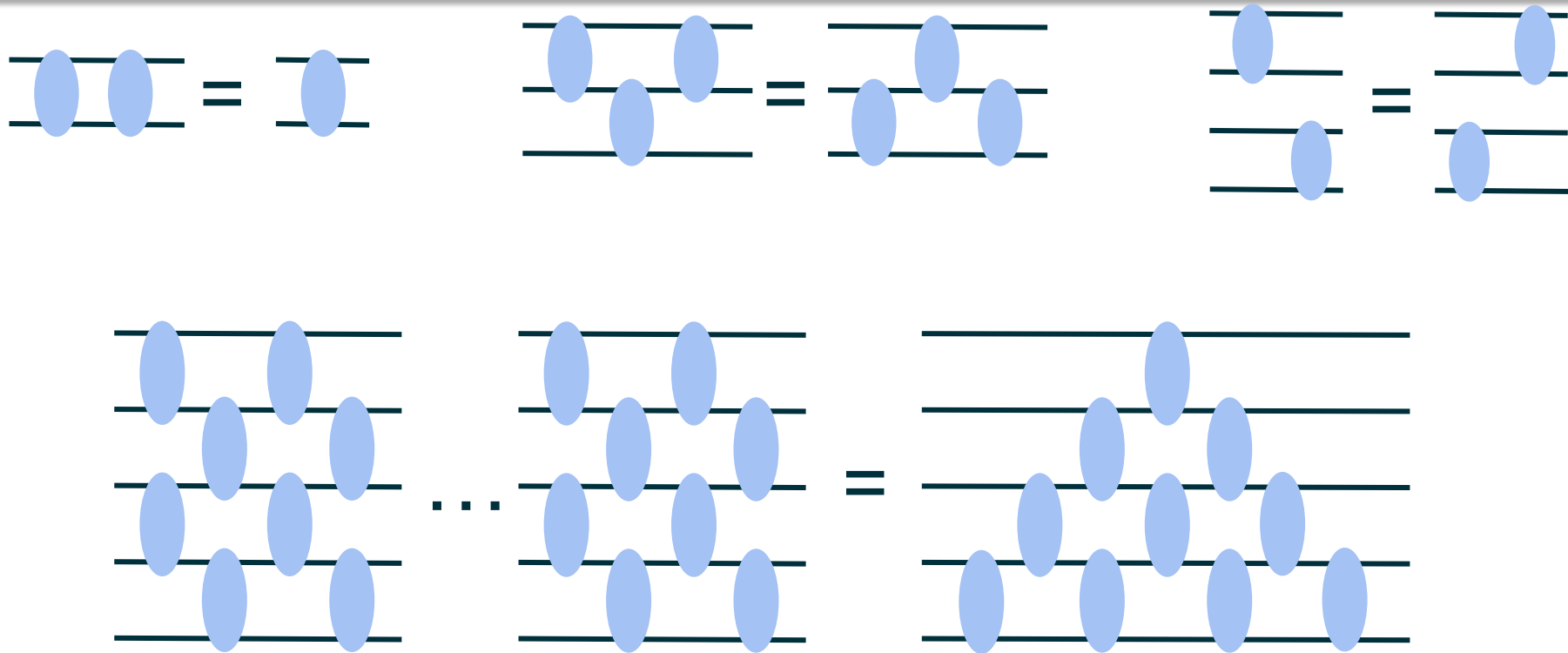
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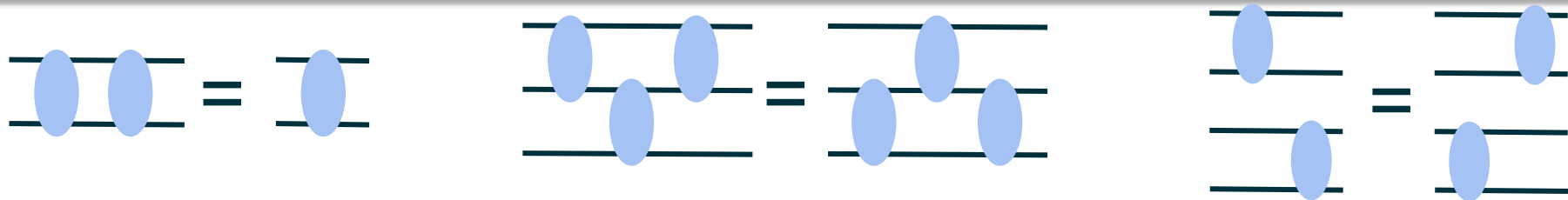
- Triangle can eat any other block!
- This allows us to reduce the number of gates significantly for certain models



# Simplification of blocks



# Free fermion models consist of blocks!



In [1-2], we show that the following models can be represented and compressed via the following blocks

[1] E. Kökcü et al (2022), PRA 105(3), 032420

[2] D. Camps et al (2022), SIAM, 43(3), 1084-1108.

## Kitaev Chain / XY model

$$2i-1 = \begin{array}{c} \text{X} \\ \text{X} \end{array}$$

$$2i = \begin{array}{c} \text{Y} \\ \text{Y} \end{array}$$

## Transverse Field Ising Model

$$2i-1 = i - \text{Z} -$$

$$2i = \begin{array}{c} \text{X} \\ \text{X} \end{array}$$

## Transverse Field XY Model

$$i = \begin{array}{c} i - \text{Z} - \text{X} - \text{Y} - \text{Z} - \\ i+1 - \text{Z} - \text{X} - \text{Y} - \text{Z} - \end{array}$$

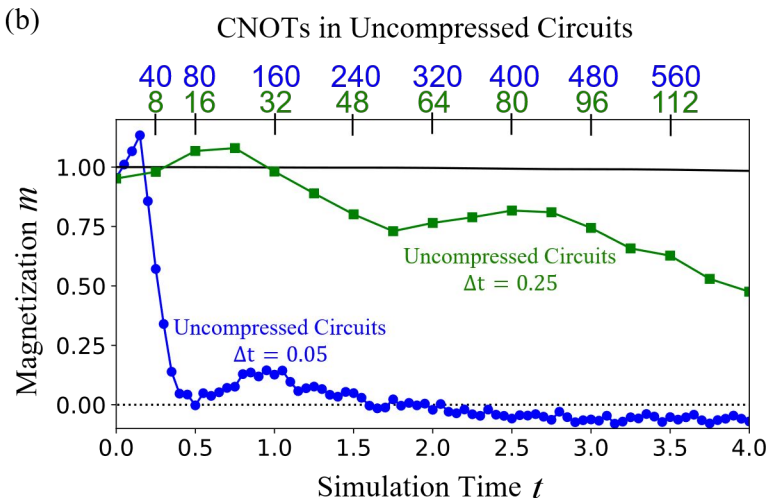
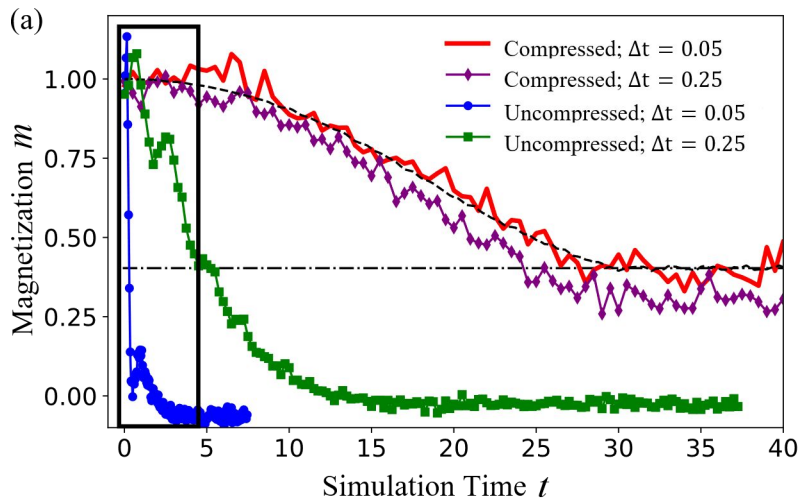


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# 5-site Transverse Field Ising

$$\mathcal{H}_{ASP}(t) = J(t) \sum_{i=1}^{n-1} X_i X_{i+1} + h_z \sum_{i=1}^n Z_i \quad \langle m(t) \rangle \equiv \frac{1}{n} \sum_i \sigma_i^z(t)$$

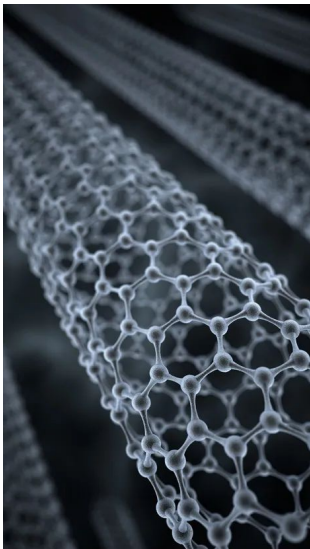
*ibmq\_brooklyn* results:



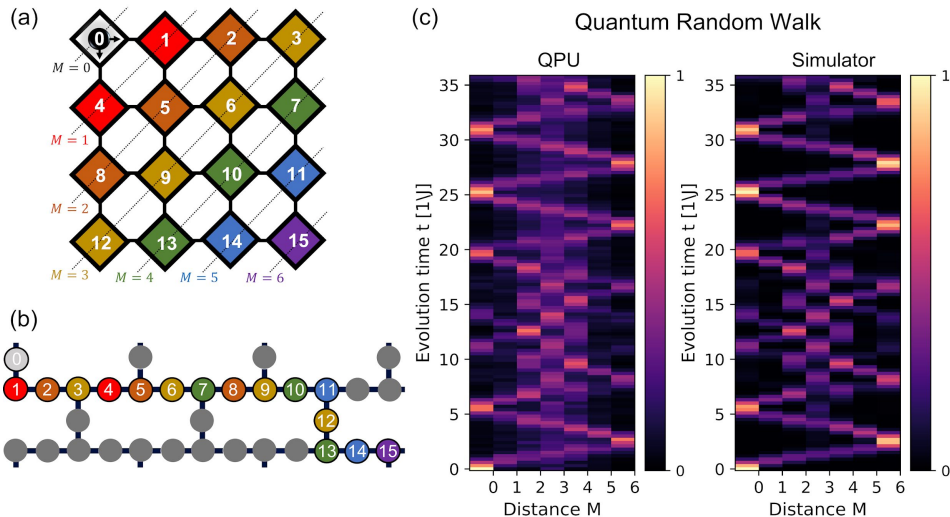
# 4x4 free tight binding model

We extend algebraic compression to free fermions with long range interactions as well, and simulate a 2-D 4x4 tight binding model on *ibmq\_washington*:

E. Kökcü et al (2023), arXiv:2303.09538



<https://www.newscientist.com/article/2093356>



# Conclusions and outlook

- We introduced a method to compress time evolution of free fermionic models on any graph
- We are applying the same method to impurity models (in progress)
- We introduced a method to compress time evolution of free fermionic models on any graph

